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Monterey, California



BUCKLING OF SHIP GRILLAGES

by

D.A. Danielson
D.P. Kihl

September 1996

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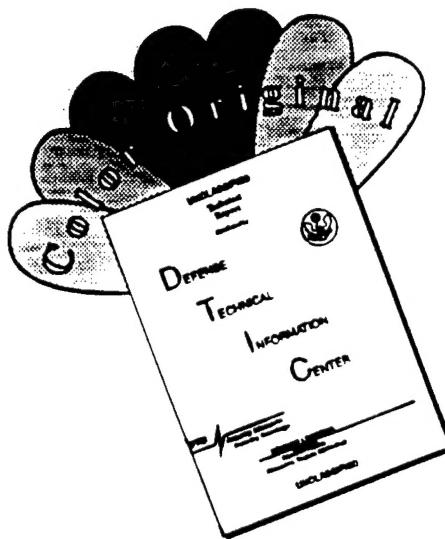
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ABSTRACT

The subject of this report is the mechanical behavior of stiffened plates, basic structural components of ships and submarines. The buckling loads of grillages subjected to axial compression with and without lateral pressure are calculated using a finite element based eigenvalue analysis. Insights are obtained into the ways in which the buckling loads and modes vary with various grillage dimensions.

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INTRODUCTION

Stiffened plates are basic structural components of ships and submarines. These structures are designed with generous safety margins against overall collapse triggered by buckling. The object of mathematical modeling is to determine design criteria to inhibit buckling at any stress less than yield. In our earlier work (see References), we have reviewed the existing literature and developed new methods for predicting the buckling loads of simple plate-beam structures. In the present work, we calculate the axial buckling loads of grillages with the use of a well-known finite element code. The goal is to determine the effect on the buckling loads of varying various grillage parameters.

The grillage now modeled consists of 3 base plates with 4 longitudinal and 2 transverse stiffeners. The dimensions of this grillage are denoted in the report by

a1 = length between transverse stiffeners

a2 = length between transverse stiffener and grillage end

b1 = width between longitudinal stiffeners

b2 = width between outer longitudinal stiffener and grillage side = b1/3

t1 = thickness of inner base plate

t2 = thickness of outer base plates = $t1 \times \frac{4}{3}$

dw1 = depth of longitudinal webs

tw1 = thickness of longitudinal webs

df1 = depth of longitudinal flanges

tf1 = thickness of longitudinal flanges

dw2 = depth of transverse webs

tw2 = thickness of transverse webs

df2 = depth of transverse flanges

tf2 = thickness of transverse flange.

The material chosen for this study is isotropic steel with Young's modulus $E = 3 \times 10^7$ psi and Poisson's ratio $\nu = .3$. The imposed boundary conditions are:

- (1) One end of the grillage has all 3 displacement components zero and all 3 rotation components zero.
- (2) The other end of the grillage where the force is applied has axial displacement constant with the other 2 displacement components zero and all 3 rotation components zero.
- (3) The ends of the transverse stiffeners have vertical displacement zero but the other 2 displacement components and all 3 rotation components are free.
- (4) All other nodes on the plates edges are completely free.

Because of the boundary conditions and the fact that the thickness t_2 of the outer base plates is chosen to be $\frac{4}{3}$ times the thickness t_1 of the inner base plate, the buckling deflections of the outer bays are usually small compared to the buckling deflections of the central bay. Thus we calculate the buckling stress in the *central* bay, defined as the initial end force divided by the cross-sectional area of the central bay, even though actual buckled members may be elsewhere.

The finite element code used is MSC Nastran/Patran. The pre-processing and post-processing is done with Patran (versions 1.4 and 1.5). The buckling analysis (solution 105) is done with Nastran (version 68). The base plates, flanges, and webs are each modeled with Quad 4 plate elements. The mesh size is denoted in this report by

m_1 = mesh length of all elements, mesh width of base plate elements

m_2 = mesh width of web and flange elements.

That is, the base plating has square elements of size $m_1 \times m_1$, while the webs and flanges have rectangular elements of size $m_1 \times m_2$. Note that this is a full finite element model of the entire grillage and does not assume symmetry. We use linear geometry for the axial load only cases, nonlinear geometry for the pressure cases, and linear material properties for all

cases. All computations are performed on a Silicon Graphics Indy Workstation.

The first phase of our work is to compare the finite element model with experimental results on a small scale grillage under axial and pressure loading. The second phase is to vary the dimensions on a full scale model to assess changes in buckling stress and mode under axial load only.

EXPERIMENTALLY TESTED GRILLAGES

We first consider the grillages (Figure 1) tested in the USNA Ship Structures Laboratory testing rig (Figure 2). We choose the following values for the model parameters:

$$a_1 = 36"$$

$$a_2 = 36"$$

$$b_1 = 9"$$

$$t_1 = \frac{3}{16}$$

$$d_{w1} = 3"$$

$$t_{w1} = .114"$$

$$d_{f1} = 1.844"$$

$$t_{f1} = .171"$$

$$d_{w2} = 9"$$

$$t_{w2} = .25"$$

$$d_{f2} = 3"$$

$$t_{f2} = .25"$$

$$m_1 = 1"$$

$$m_2 = .5".$$

We find that the buckling load is not significantly changed by adopting smaller values for the mesh sizes m_1 or m_2 .

The prebuckling stage is approximately one of uniform axial compression with little bending (Figure 3). The grillage buckles under a force of 470 kips. The buckling mode is a deformation involving primarily bending with little stretching (Figure 4). The inner base plate buckles into a square quilt of half-wavelength 9" (Figure 5). The webs (both transverse and longitudinal) also buckle with half-wavelength 9" (Figures 6-7). The flanges mainly just twist about their center line. We call this buckling mode TRIPPING to be consistent with

our earlier papers. *

We also ran nonlinear analyses (solution 106) for the cases of bottom pressure $p = 0$, $p = 5$ psi, or $p = 20$ psi. For the case $p = 0$, the plate only compresses in-plane and does not buckle. However, the buckles are clearly visible by the time the axial load reaches 475 kips for the $p = 5$ psi case and 450 kips for the $p = 20$ psi case. But since the buckling pattern develops gradually as the axial load increases, it is hard to pin down an exact "buckling" load. It is interesting that the middle plate develops only 3 half waves (instead of 4 as in the buckling analysis for $p = 0$). The magnitude of displacement for the $p = 20$ psi case is shown in Figure 8 with no axial load, and in Figure 9 with 550 kips axial load. If we increase the load much further than 500 - 550 kips, the displacements become large and the code will not give a solution (Figure 10).

A table comparing model results to experimentally measured collapse loads for 6 nominally identical grillages is given below:

p (psi)	Model (kips)	Experiments (kips)
0	470	326, 312, and 301
5	less than 475	316
10	not calculated	306
20	less than 450	296

*In our earlier work, we did not allow the stiffener cross-section to deform in its plane. This assumption led to a predicted buckling load which was much too high.

Note that both theory and experiments indicate that the presence of lateral pressure (on the plating side) tends to weaken the grillages. The disparity between the buckling loads of the perfect model and the experimental measurements is probably due to geometric imperfections, residual stress, differing boundary conditions, nonlinear material properties, etc. in the tested grillages.

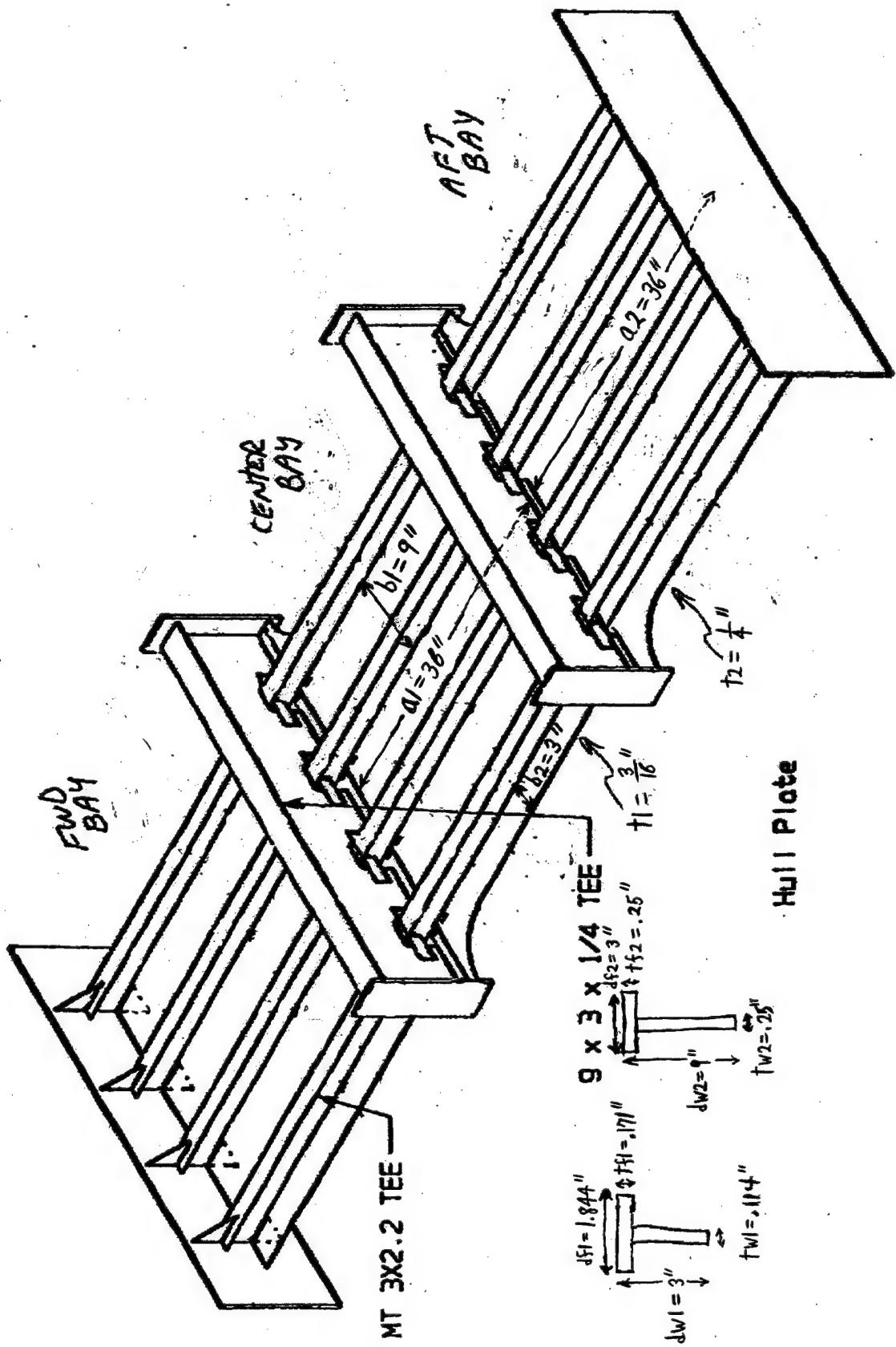


Figure 1

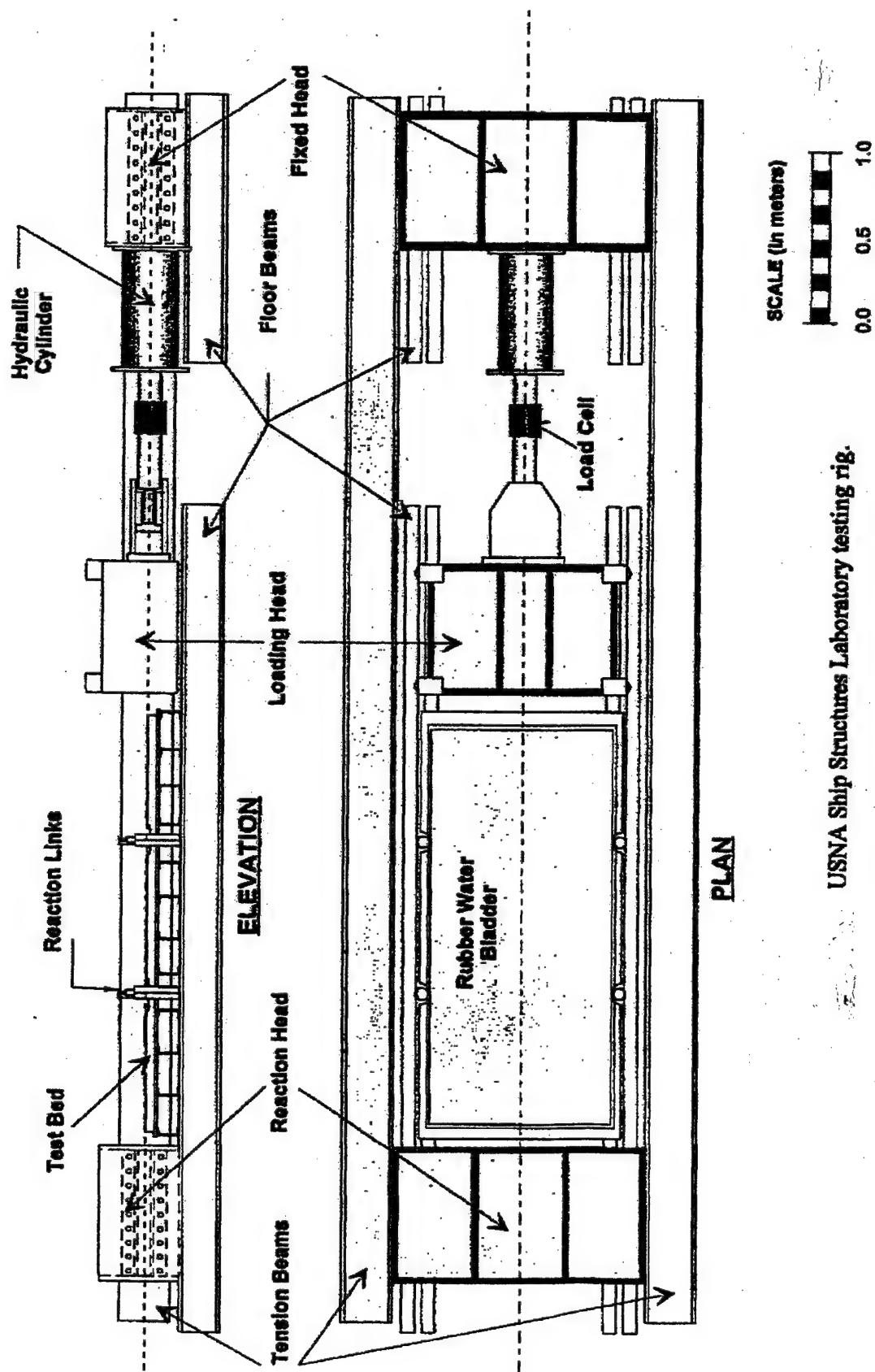
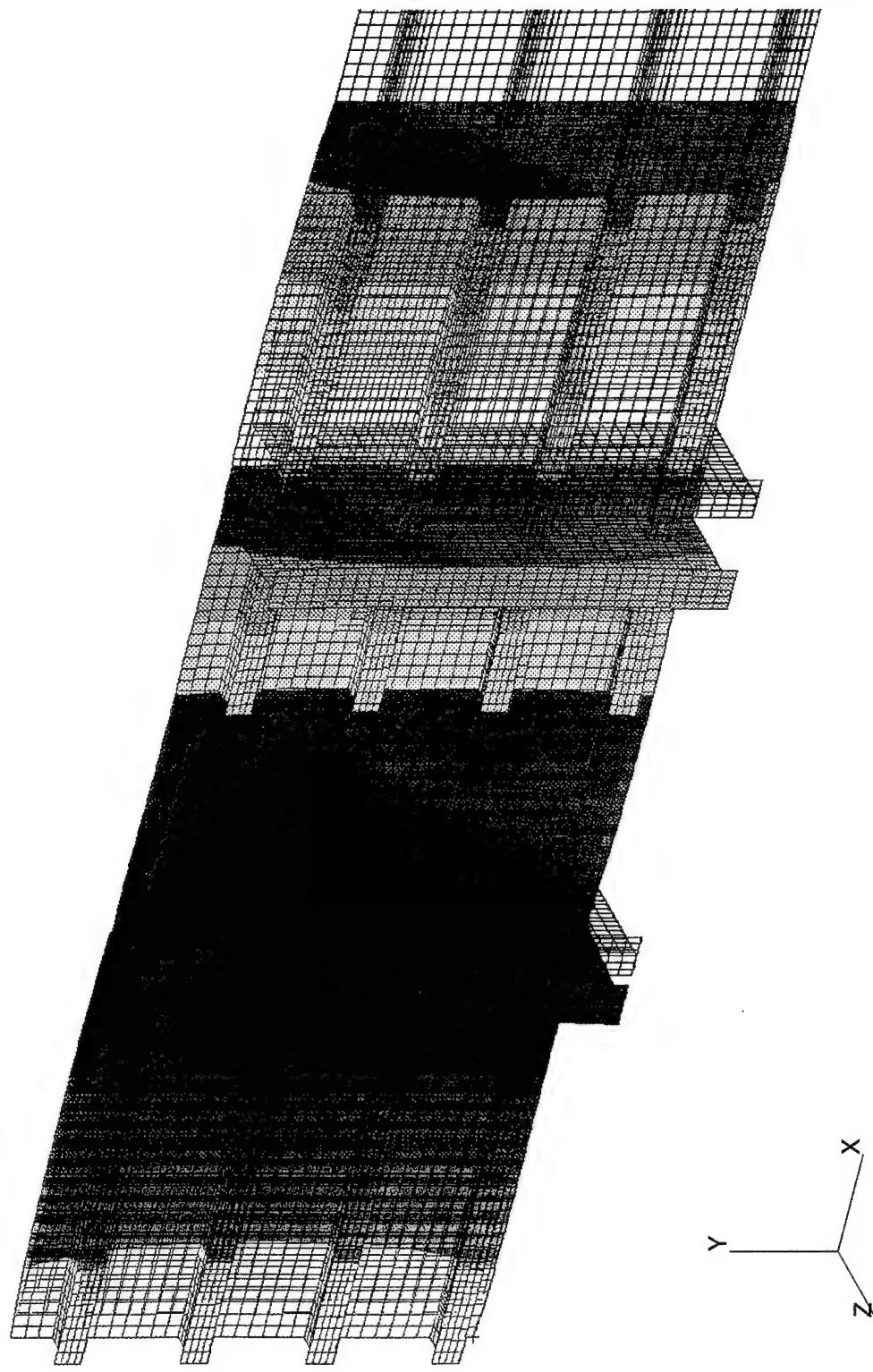


Figure 2

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DEFORMATION PLOT LC=1.23 RES=2.1 MSC/PATRAN R-1.4 MSC/NASTRAN 29-Jan-96 11:55:35

Prebuckling magnitude of displacements



.000000000000

.00000002400

.00000004700

.00000007100

.00000009400

.0000001180

.0000001420

.0000001650

.0000001890

.0000002130

.0000002360

.0000002600

.0000002830

.0000003070

.0000003310

.0000003540

Figure 3

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DEFORMATION PLOT LC=2.24 RES=3.1 MSC/PATRAN R-1.4 MSC/NASTRAN 29-Jan-96 13:17:22

470 kips - Magnitude of displacement

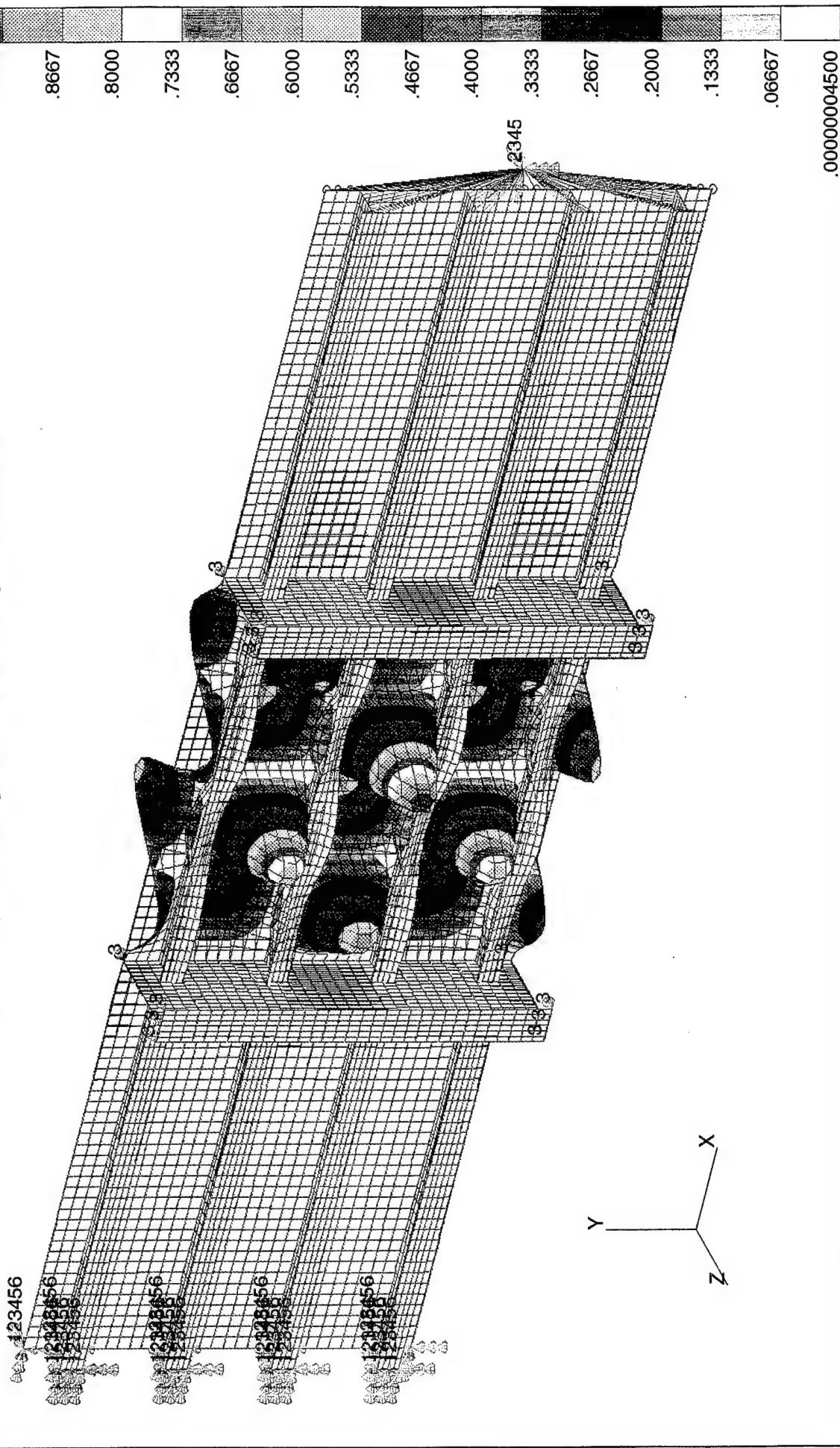


Figure 4

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DEFORMATION PLOT LC=2.24 RES=3.1 MSC/PATRAN R-1.4 MSC/NASTRAN 29-Jan-96 13:31:39

470 kips - z displacement

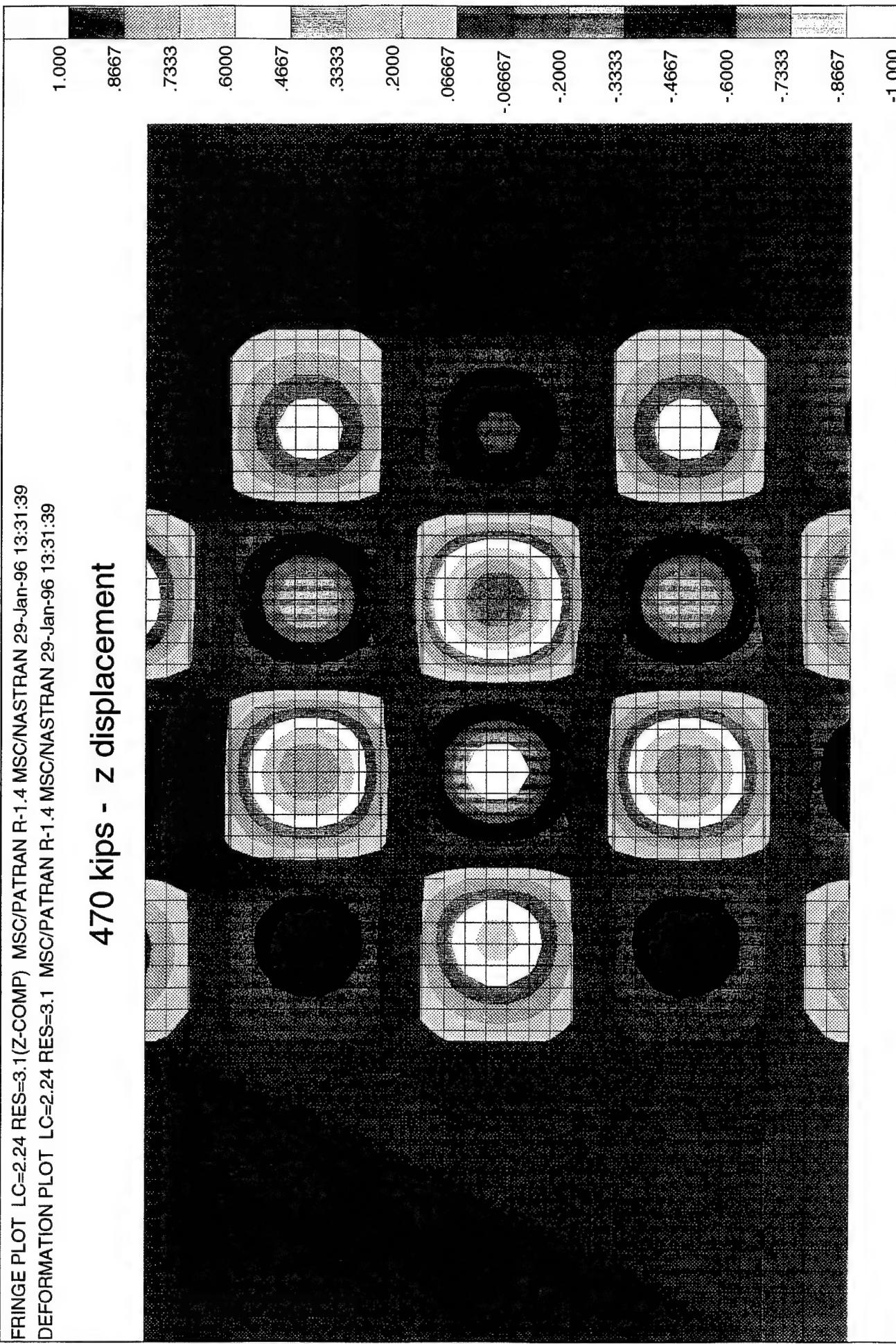


Figure 5

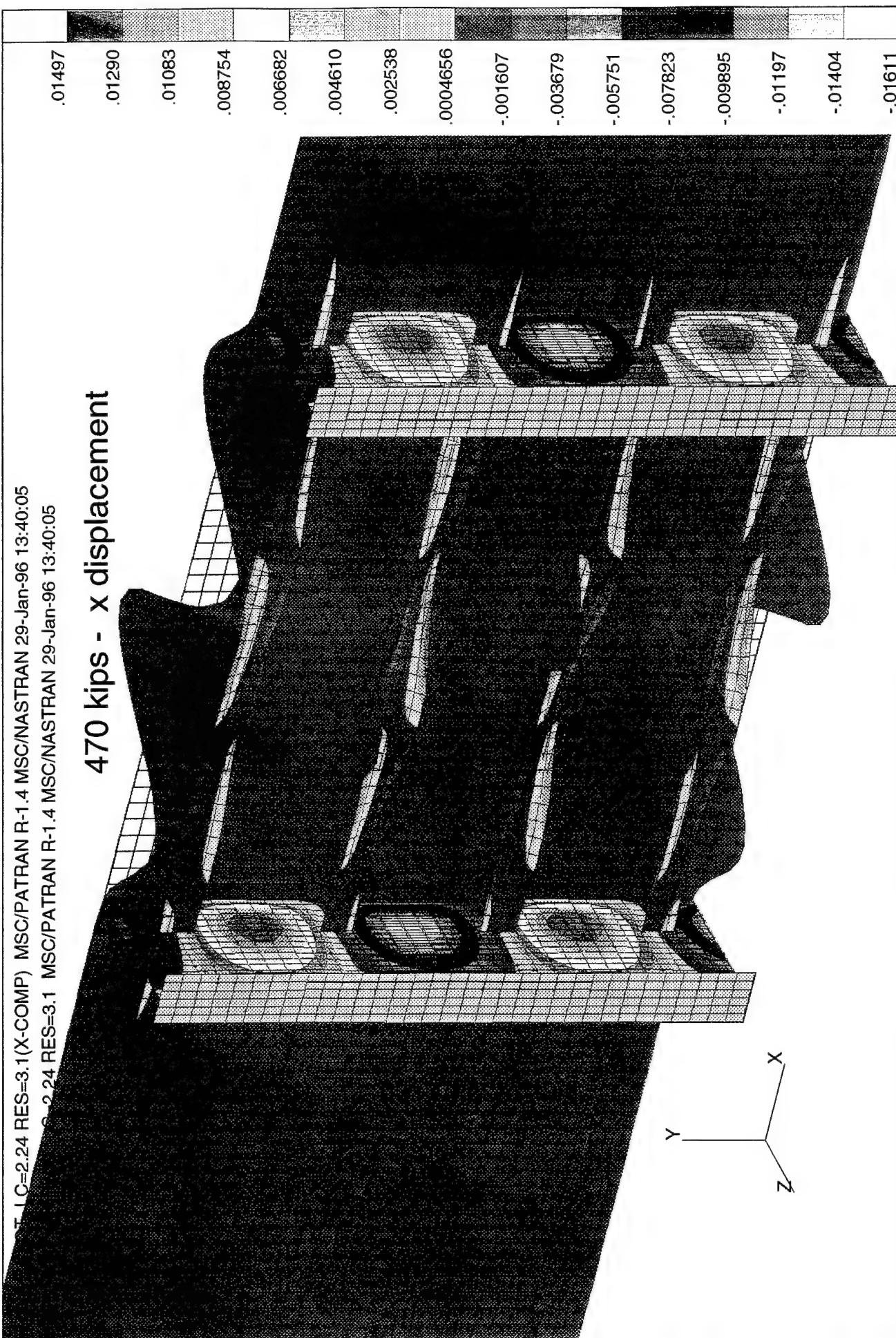


Figure 6

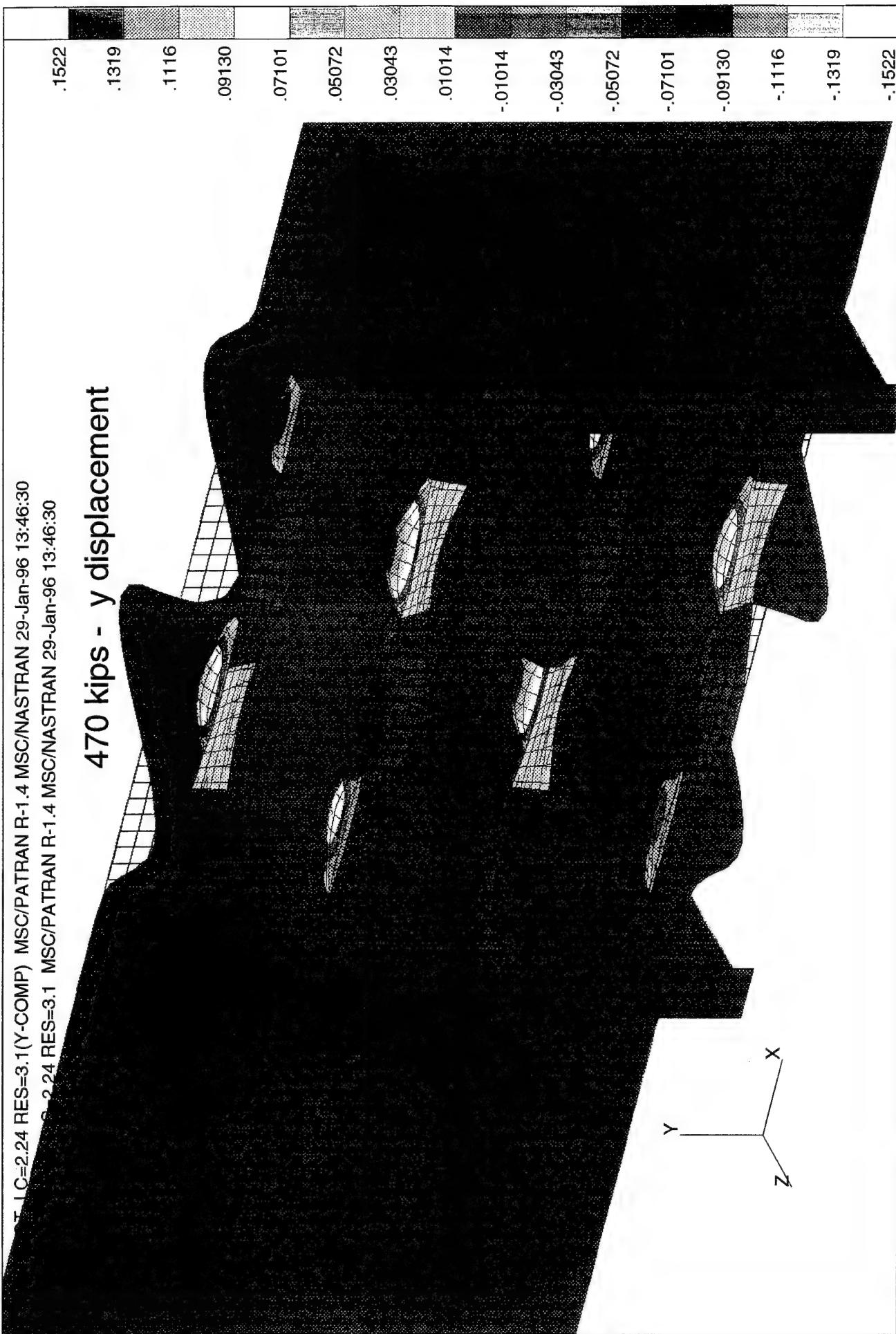


Figure 7

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DEFORMATION PLOT LC=4.79 RES=2.1 MSC/PATRAN R-1.4 MSC/NASTRAN 12-Feb-96 10:18:49

$p=20$ psi - Magnitude of displacement

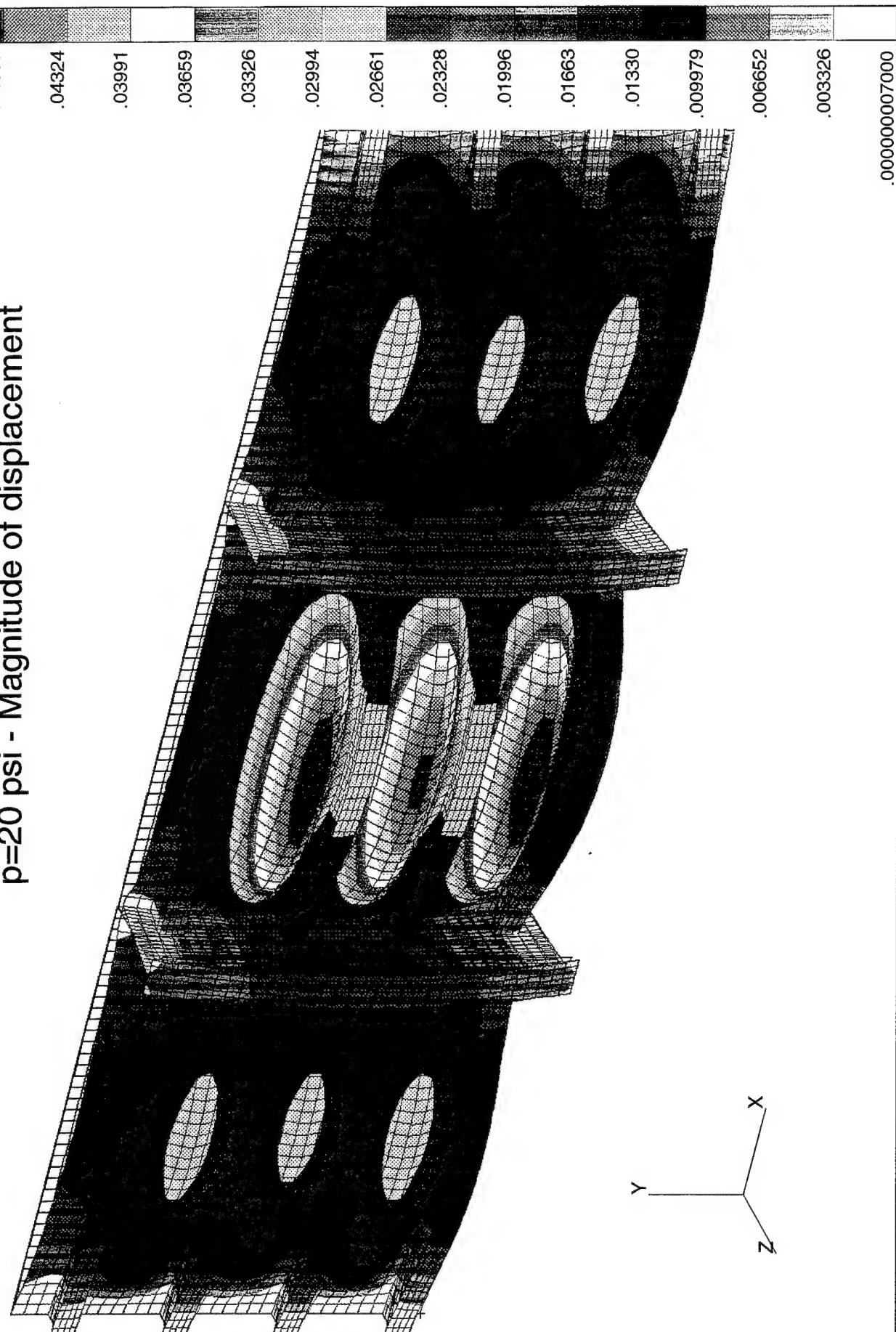


Figure 8

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DEFORMATION PLOT LC=3.90 RES=2.1 MSC/PATRAN R-1.4 MSC/NASTRAN 12-Feb-96 10:37:53

$p=20$ psi, $F=550$ kips - Magnitude of displacement

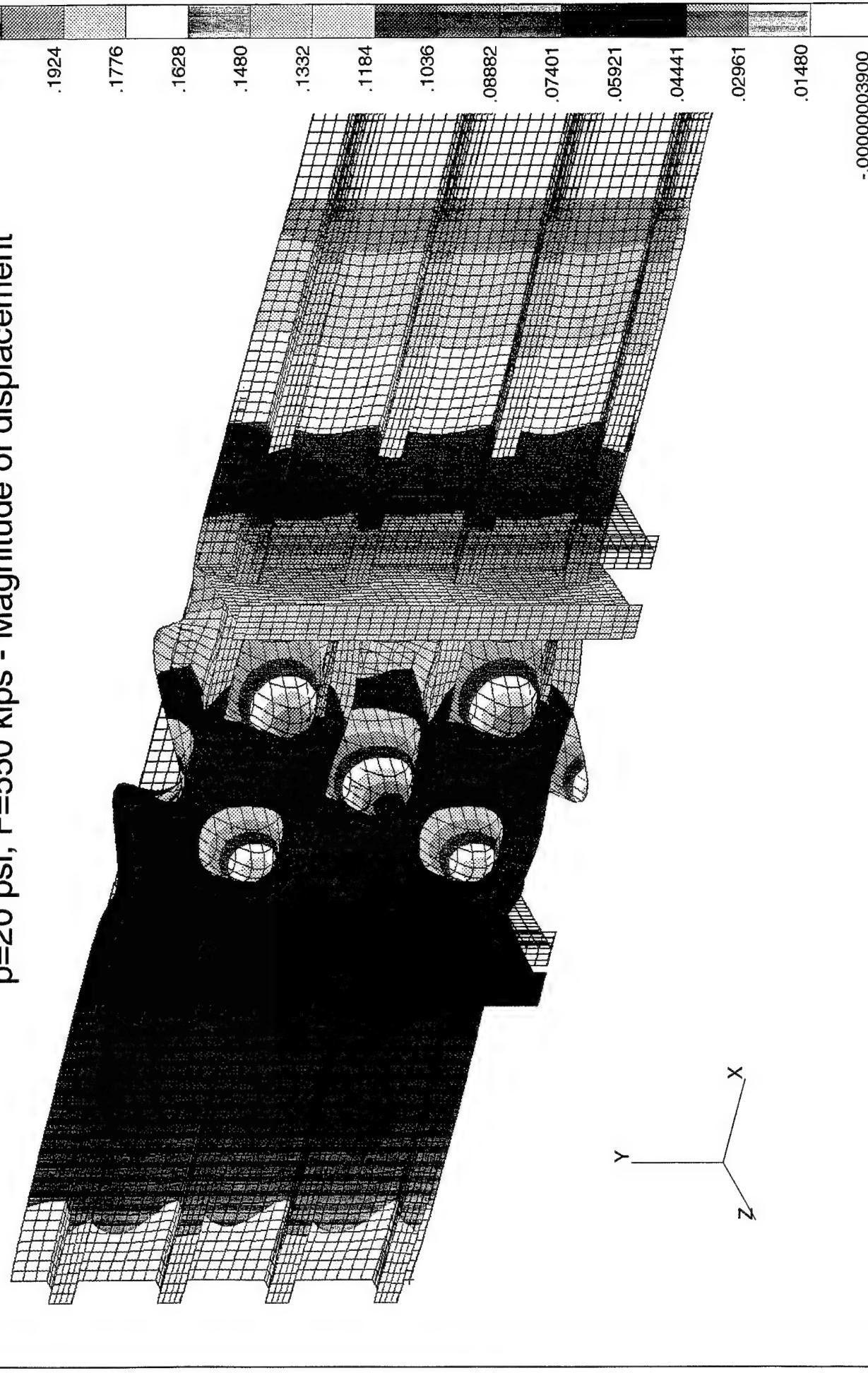


Figure 9

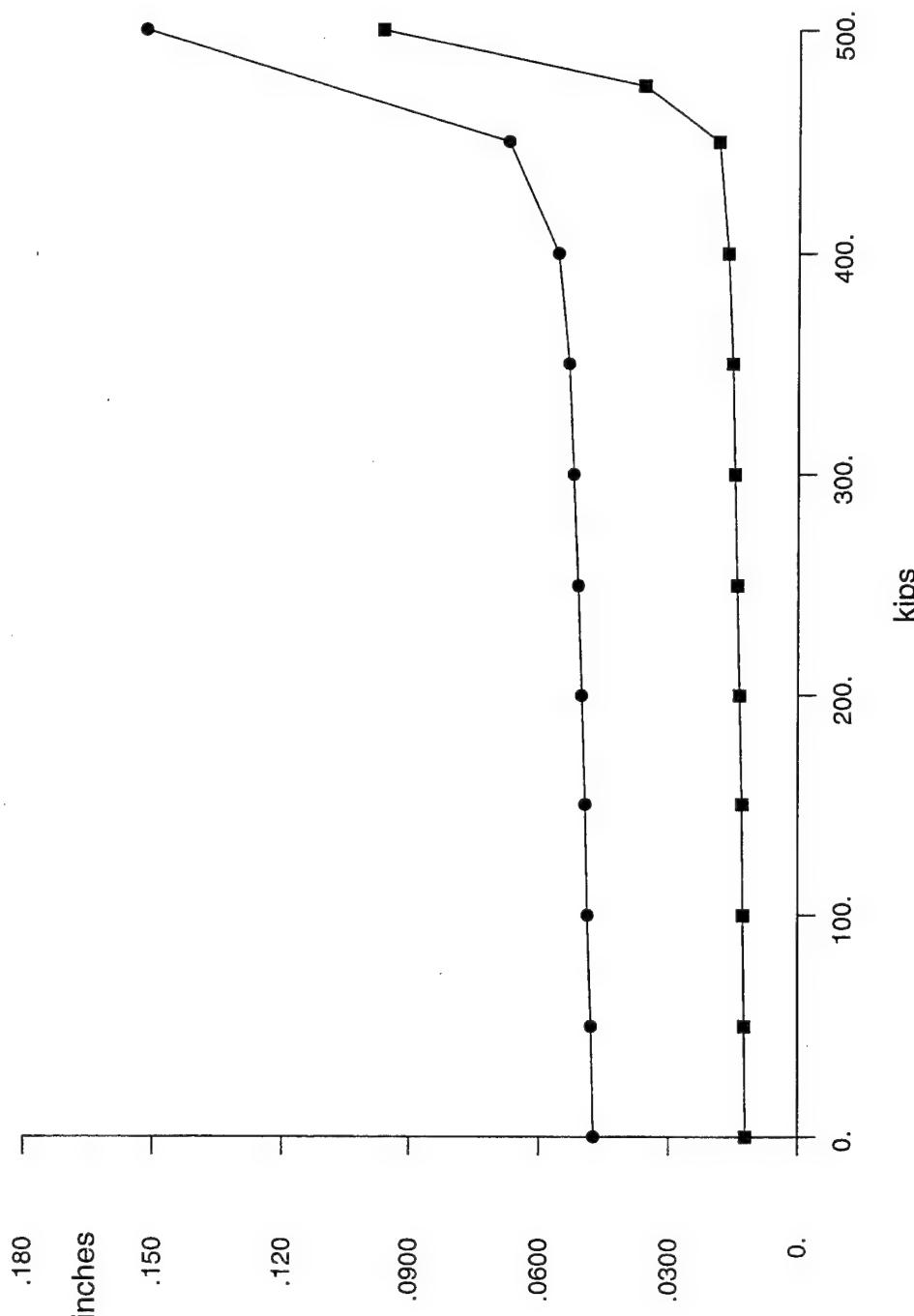
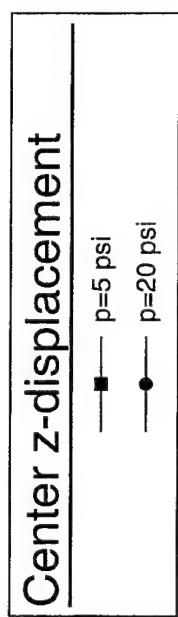


Figure 10

GRILLAGES OF VARIOUS DIMENSIONS

Generic model

We define the generic model as a grillage with the following parameter values:

$$a1 = 96''$$

$$a2 = 48''$$

$$b1 = 24''$$

$$t1 = \frac{9}{16}''$$

$$dw1 = 6''$$

$$tw1 = .2''$$

$$df1 = 4''$$

$$tf1 = .225''$$

$$dw2 = 12''$$

$$tw2 = .2''$$

$$df2 = 4''$$

$$tf2 = .225''$$

$$m1 = 2''$$

$$m2 = 1''$$

The buckling stress in the central bay of the generic model is 66.2 ksi. The buckling mode is TRIPPING of the central bay (Figure 11). The half-wavelength of the buckles is 24".

Note that the outer bays of the generic model are made half as long as the inner bay, to reduce the size of the finite element model. The buckling load of the full scale model differs from that of the generic model by an insignificant amount.

For the remainder of the report, we consider various grillages with only one dimension different from those of the generic model.

Variation of a1, length between transverse stiffeners

We consider 6 grillages having differing values of a1. All of the other parameters are the same as for the generic grillage, except for the case a1 = 144" we choose the mesh length m1 = 4" to make the program size smaller. The buckling stress decreases monotonically with a1 (Figure 12). The buckling mode is TRIPPING in all cases. For the 5 cases with a1 a multiple of 24", the half-wavelength of the buckles is 24". For the case a1 = 60", the half-wavelength is 30" (Figure 13).

Variation of b1, width between longitudinal stiffeners

We consider 4 grillages having b1 differing multiples of 12". The buckling stress decreases monotonically with b1 (Figure 14). The buckling stress is somewhat larger than the classical buckling stress of a square simply-supported plate (Figure 15):

$$\frac{E\pi^2(t1)^2}{3(1 - \nu^2)(b1)^2}$$

The buckling mode is TRIPPING (similar to Figure 11) and the half-wavelengths of the buckles is b1.

Variation of t_1 , thickness of inner base plate

We consider 5 grillages having differing values of t_1 . The buckling stress reaches a maximum value for an optimal t_1 (Figures 16 - 17). For the 3 lowest cases, the buckling mode is TRIPPING, with the half-wavelengths being 19.2" for the case $t_1 = \frac{3}{16}$ " and 24" for the cases $t_1 = \frac{3}{8}$ " and $t_1 = \frac{9}{16}$ ". For the cases $t_1 = \frac{3}{4}$ " and $t_1 = \frac{15}{16}$ " the buckling mode is GLOBAL. The entire central bay buckles in (Figure 18), while the stiffeners deflect inwards and laterally (Figure 19).

Variation of dw_1 , depth of longitudinal webs

We consider 7 grillages having differing values of dw_1 . The buckling stress reaches a maximum value for an optimal dw_1 (Figure 20). For the cases $dw_1 = 3$ " and $dw_1 = 4$ ", the buckling mode is GLOBAL (Figure 21). For the cases $dw_1 = 5$ ", $dw_1 = 6$ ", and $dw_1 = 9$ ", the buckling mode is TRIPPING with half-wavelength of 24". For the cases $dw_1 = 10$ " and $dw_1 = 11$ ", the buckling mode is LOCAL (Figure 22). The web and flange of a middle stiffener of the central bay buckle with a half-wavelength dw_1 . We find that the critical load for the local buckling case is not significantly changed by adopting smaller values for the web and flange mesh element width m_2 .

Variation of tw_1 , thickness of longitudinal webs

We consider 7 grillages having differing values of tw_1 . The buckling stress increases monotonically with tw_1 (Figure 23). For the case $tw_1 = .05$ ", the web of a middle stiffener of the forward bay buckles LOCALLY, and for the case $tw_1 = .1$ ", the web of a middle stiffener of the central bay buckles LOCALLY (Figure 24). For the cases $tw_1 = .15$ ", $.2$ ", $.3$ ", and $.7$ " the buckling mode is TRIPPING with half-wavelength 24" (Figure 25). For the case $tw_1 = 1$ ", the buckling mode is TRIPPING with half-wavelength 19.2".

Variation of df1, depth of longitudinal flanges

We consider 4 grillages having differing values of df1. For the cases $df1 = 2"$, $4"$, and $6"$, the buckling mode is TRIPPING and the buckling stress varies little (Figure 26). For the case $df1 = 0$ (no flange), the web of a middle stiffener of the central bay buckles LOCALLY (Figure 27).

Variation of tf1, thickness of longitudinal flanges

We consider 4 grillages having different values of tf1. For the cases $tf1 = .225"$, $.5"$, and $1"$, the buckling mode is TRIPPING and the buckling stress varies little (Figure 28). For the case $tf1 = .1"$, the flange of a middle stiffener of the central bay buckles LOCALLY (Figure 29).

Variation of dw2, depth of transverse webs

We consider 3 grillages having differing values of dw2. For all 3 cases $dw2 = 4"$, $12"$, and $16"$, the buckling mode is TRIPPING and the buckling stress varies little (Figure 30).

Variation of tw2, thickness of transverse webs

We consider 3 grillages having differing values of tw2. For the cases $tw2 = .15"$ and $.2"$, the buckling mode is TRIPPING and the buckling stress varies little (Figure 31). For the case $tw2 = .1"$, a web buckles LOCALLY under a *tensile* force (Figure 32). Compression in the web arises from Poisson ratio effects.

Variation of df2, depth of transverse flanges

We consider 2 grillages having differing values of df2. For both cases $df2 = 0$ (no flange) and $df2 = 4"$, the buckling mode is TRIPPING and the buckling stress varies little (Figure 33). Note that eliminating the transverse flange would not reduce the buckling strength of the grillage under axial compression.

Variation of tf_2 , thickness of transverse flanges

We consider 2 grillages having differing values of tf_2 . For both cases $tf_2 = .1''$ and $tf_2 = .225''$, the buckling mode is TRIPPING and the buckling stress varies little (Figure 34).

Summary

Through study of all these different results, we can gain insights into the ways in which the buckling loads and modes vary with the various dimensions.

The buckling stress decreases less with length a_1 than with width b_1 between stiffeners. The buckling stress is larger for an optimal value of plating thickness t_1 and depth dw_1 of longitudinal webs. The buckling stress increases with thickness tw_1 of longitudinal webs. The buckling stress is not much affected by increases in the dimension df_1 and tf_1 of the longitudinal flanges, dimensions dw_2 and tw_2 of transverse webs, or dimensions df_2 and tf_2 of transverse flanges. However, the buckling stress is smaller for too small values of thickness tw_1 of longitudinal webs, depth df_1 and thickness tf_1 of longitudinal flanges, and thickness tw_2 of transverse webs.

We have found three distinct failure modes which we have named TRIPPING, GLOBAL, and LOCAL. For the majority of our cases, the failure mode is TRIPPING with the inner base plate buckling into a square quilt, which forces the webs to buckle into a rectangular quilt with the same wavelength, while the flanges twist (Figure 11). However, if a square pattern cannot fit into the base plate (Figure 13), or if the relative strength of the stiffeners is much greater than the base plate (Figure 25), the inner base plate may buckle into a non-square quilt, and the flanges may also undergo lateral deformation. If the relative strength of the inner base plate is much greater than the stiffeners, the failure mode may be GLOBAL with the entire central bay buckling in (Figures 18-19 and 21). If the relative strength of a stiffener is too small, it may collapse LOCALLY before any of the other components buckle (Figures 22, 24, 27, 29, and 32).

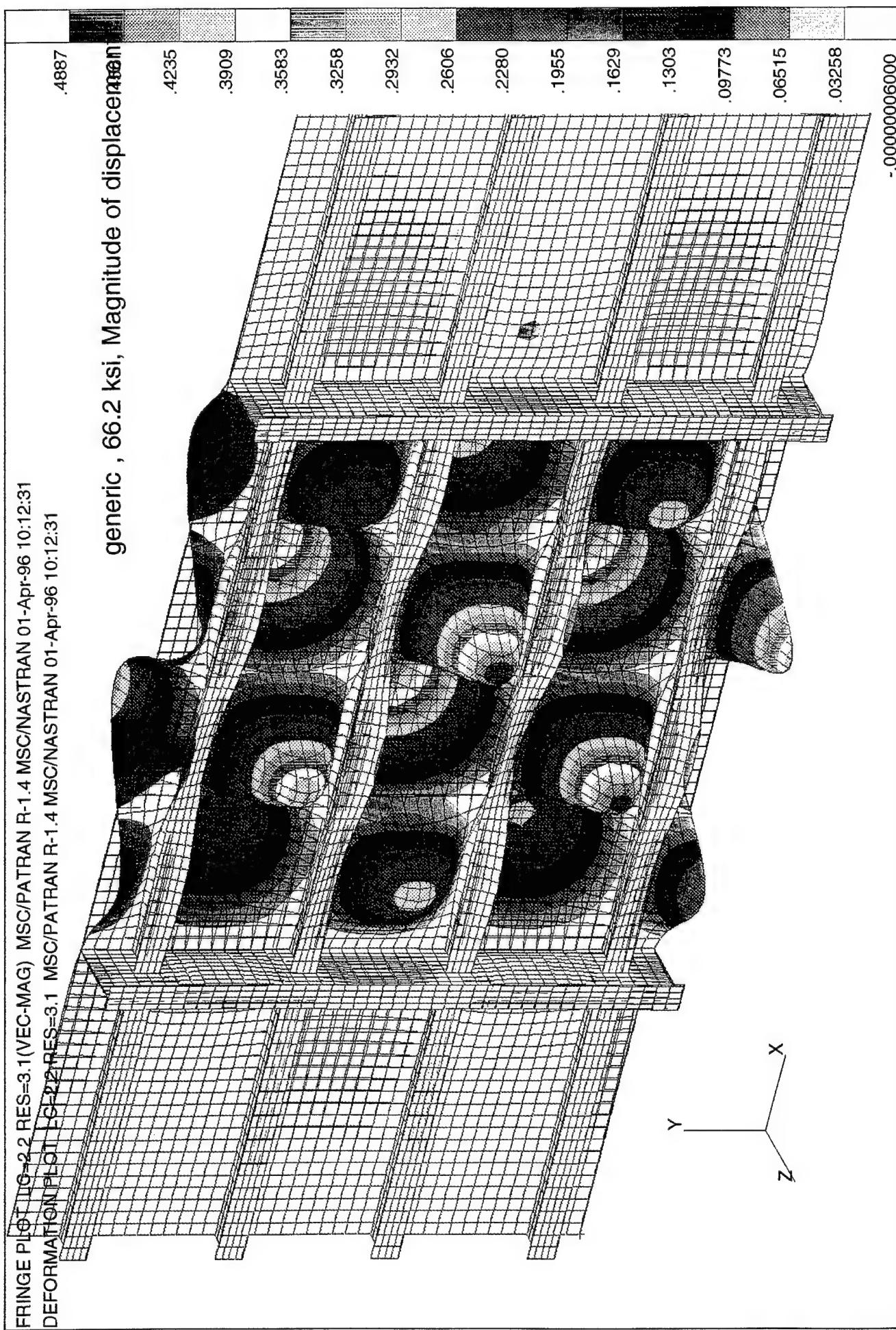


Figure 11

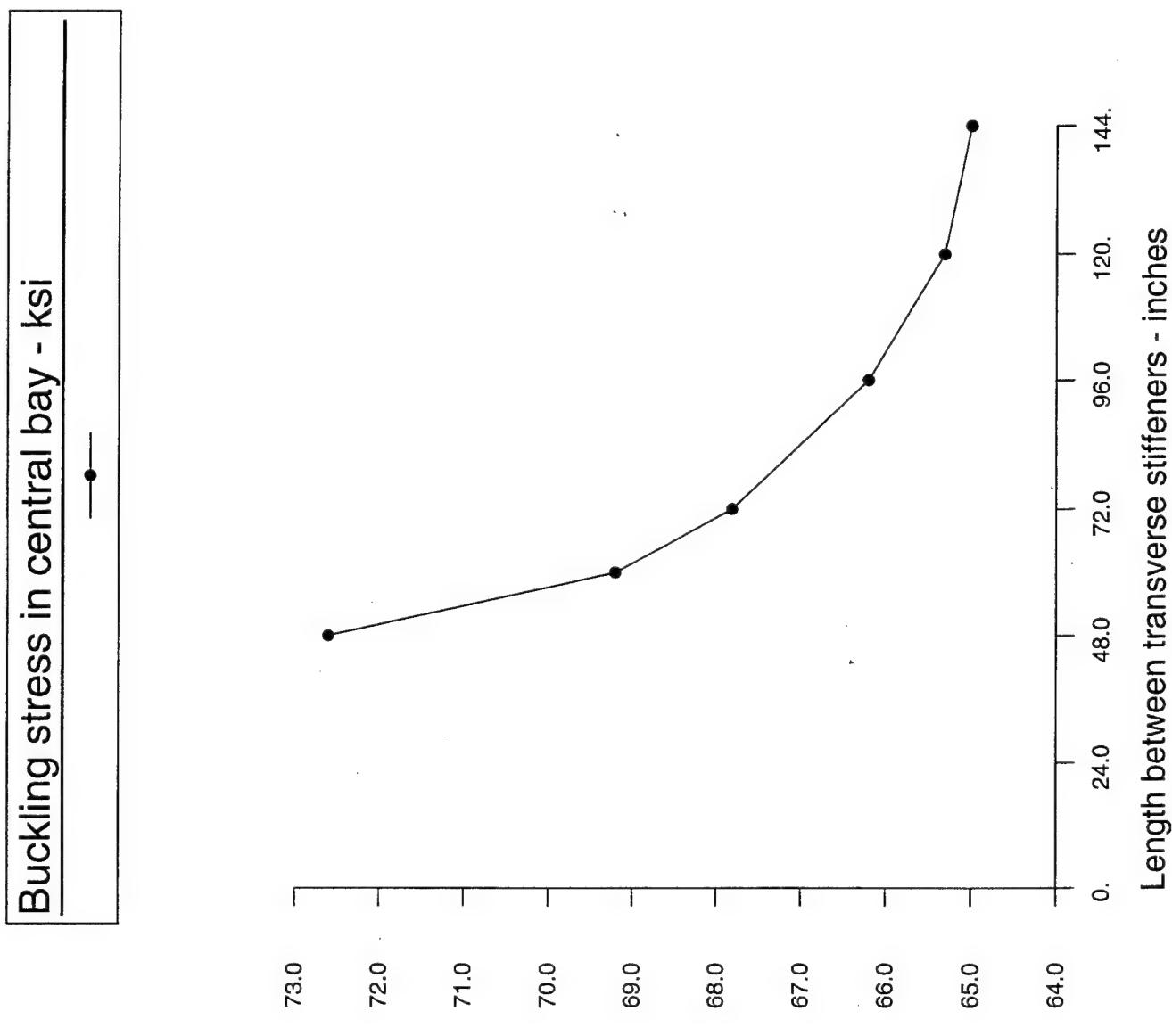


Figure 12

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a1=60", 69.2 ksi - Magnitude of displacement

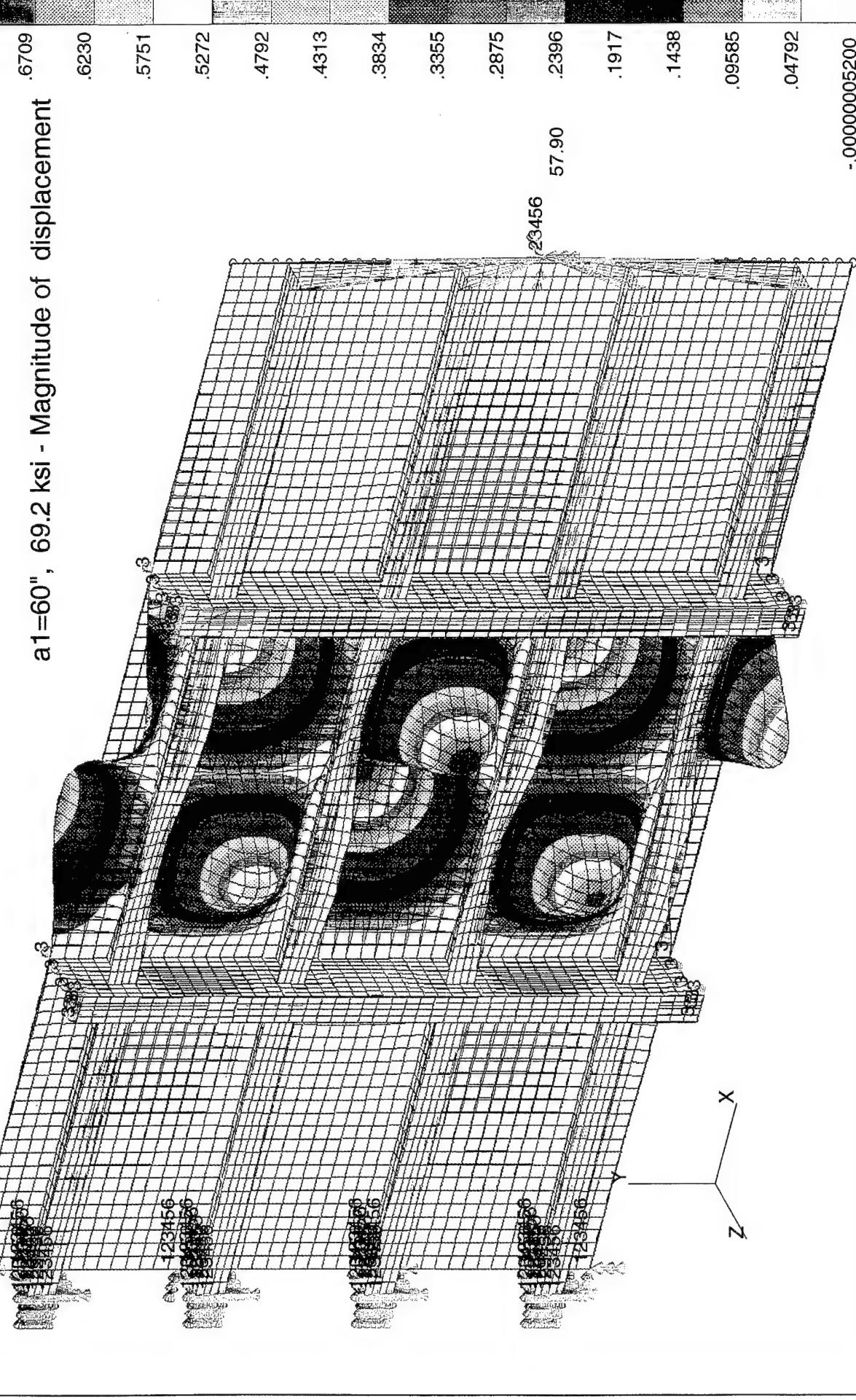


Figure 13

Buckling stress in central bay - ksi

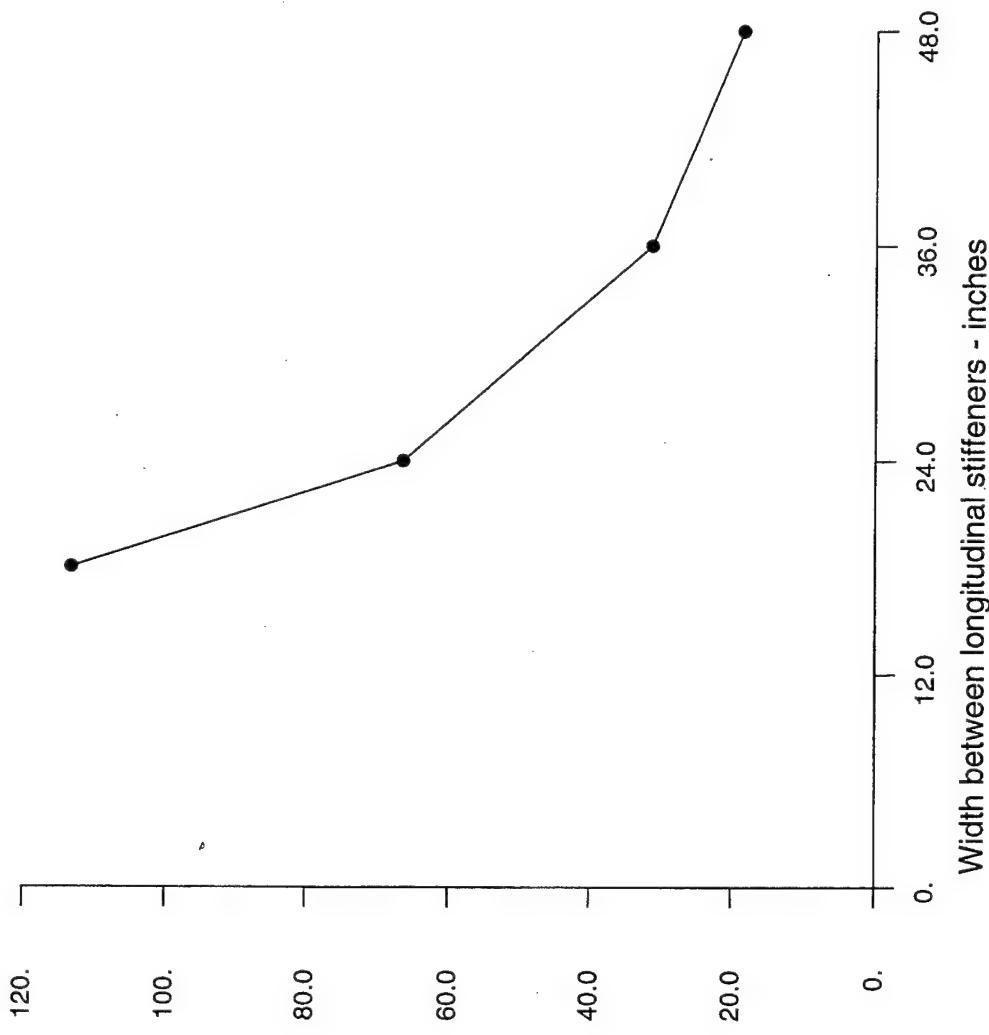


Figure 14

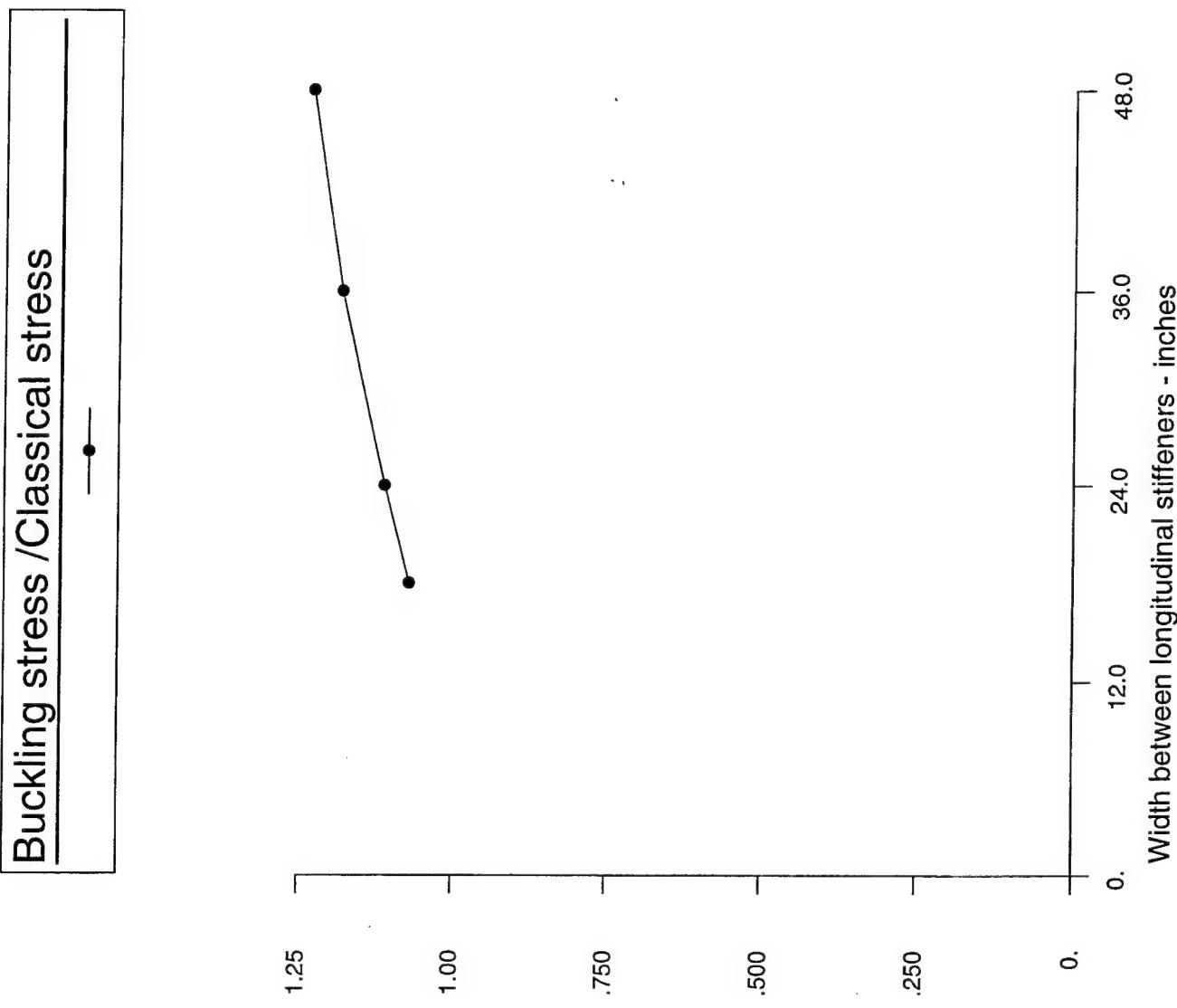


Figure 15

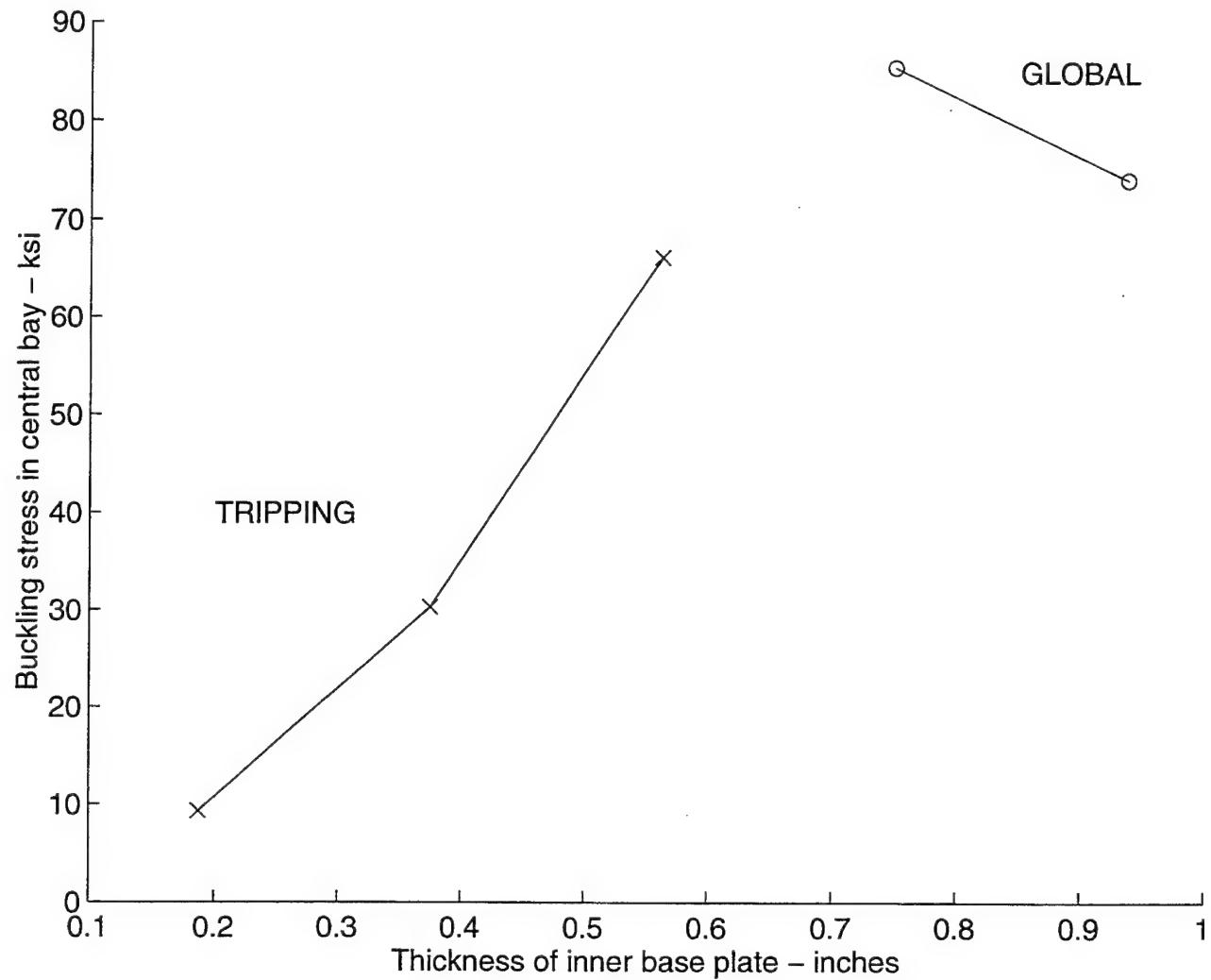


Figure 16

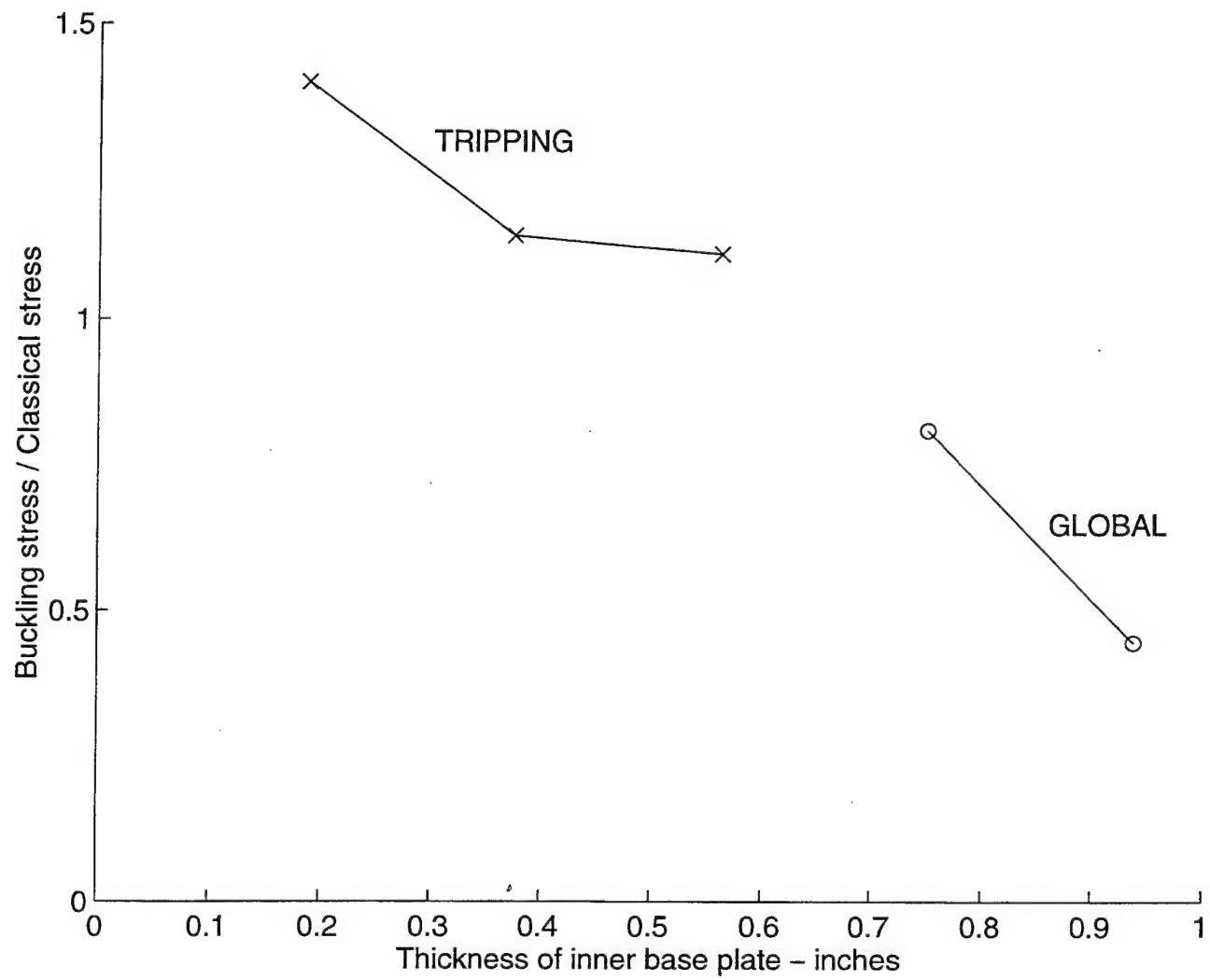


Figure 17

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t1=.9375", 74.1 ksi, z-displacement

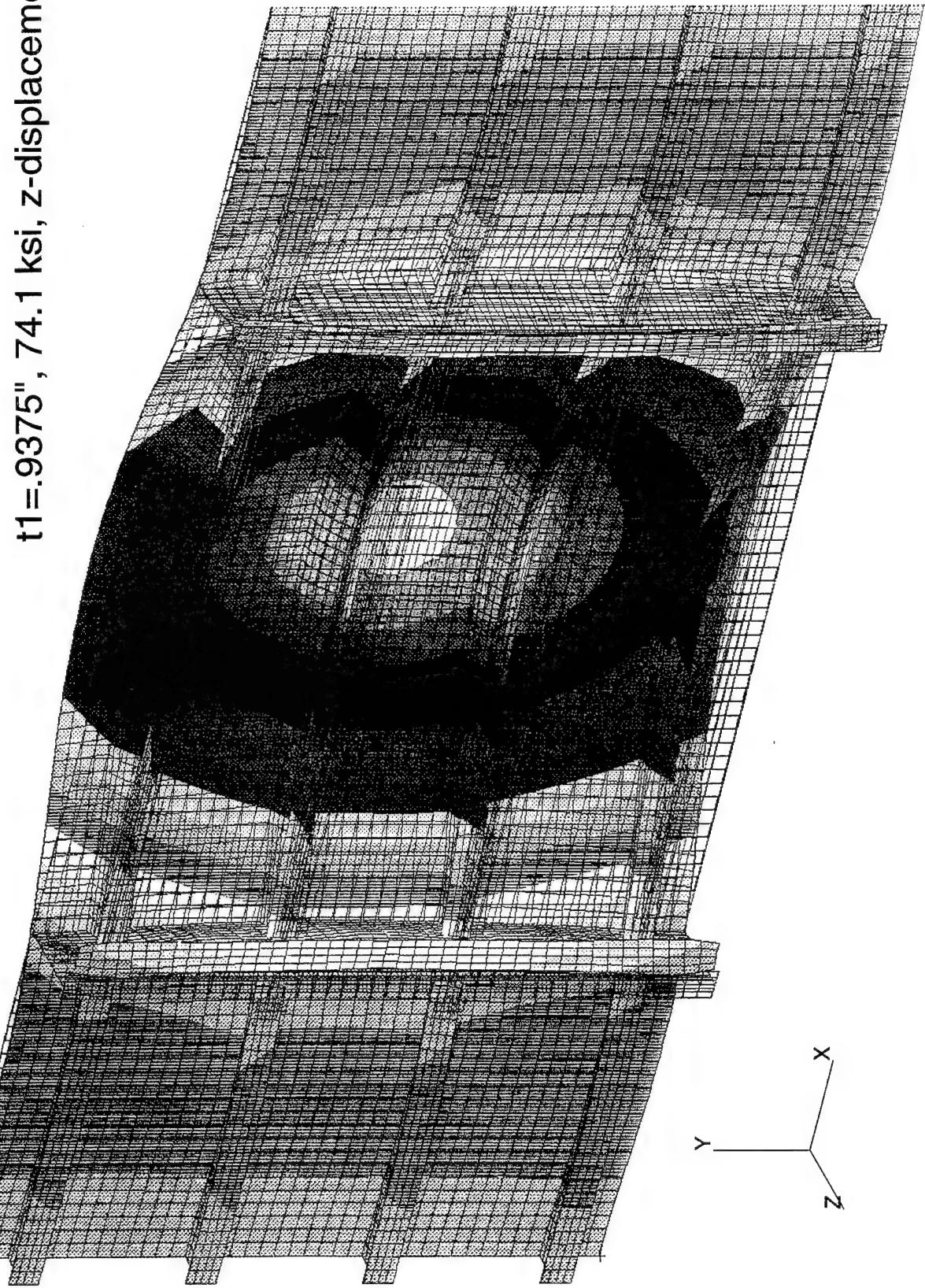


Figure 18

FRINGE PLOT LC=2.2 RES=3.1(Y-COMP) MSC/PATRAN R-1.4 MSC/NASTRAN 18-Mar-96 13:24:48
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t1=.9375", 74.1 ksi, y-displacement

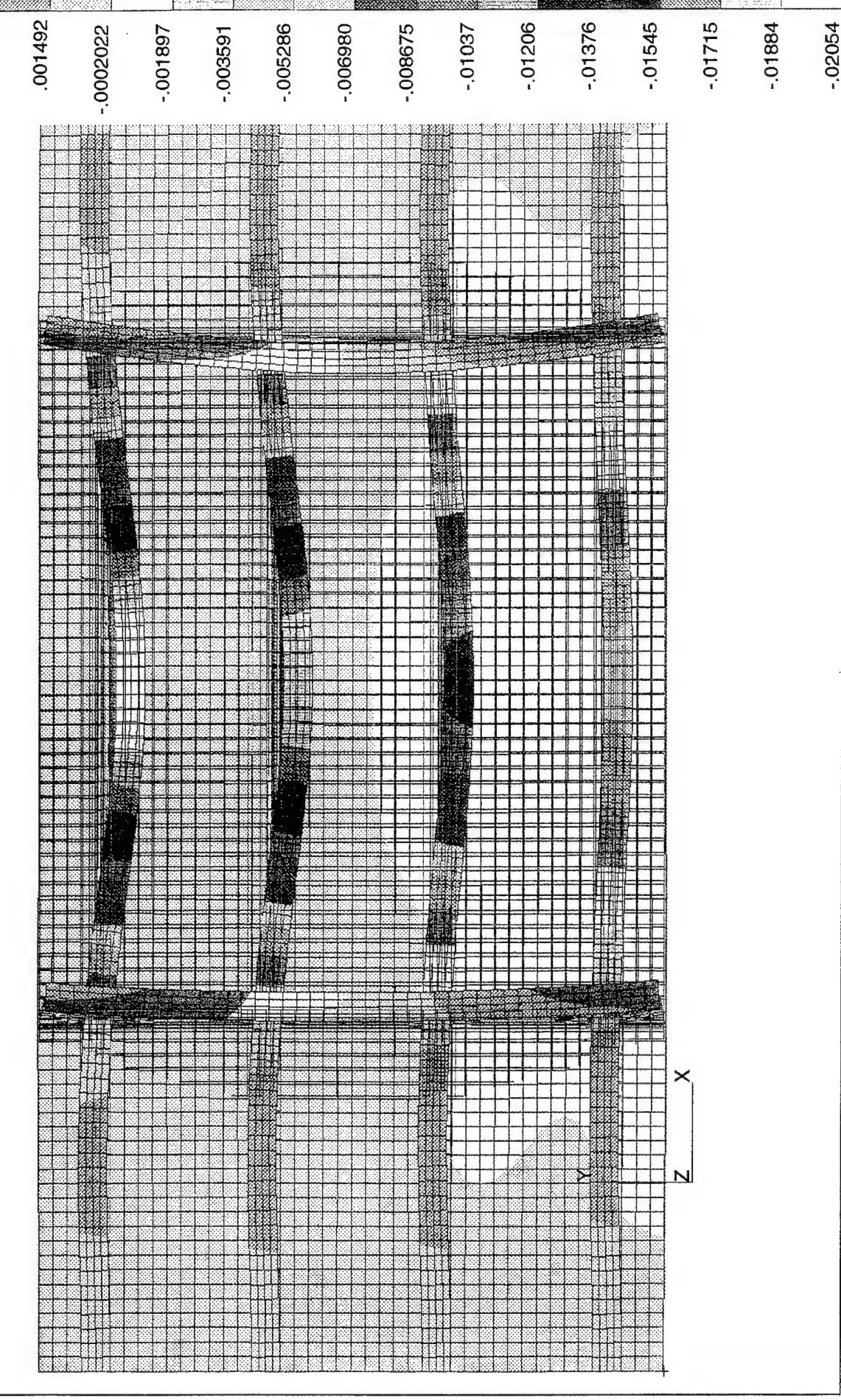


Figure 19

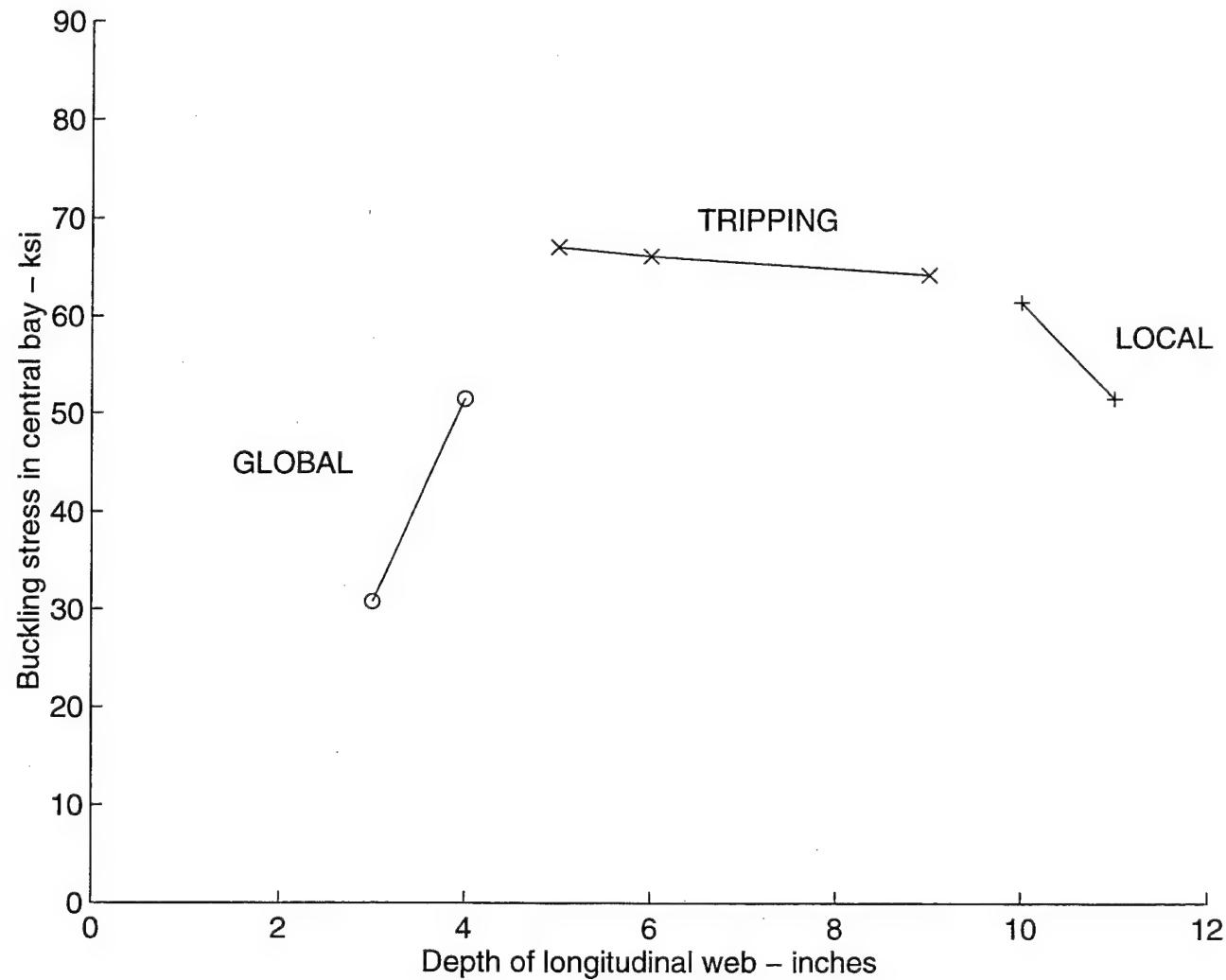


Figure 20

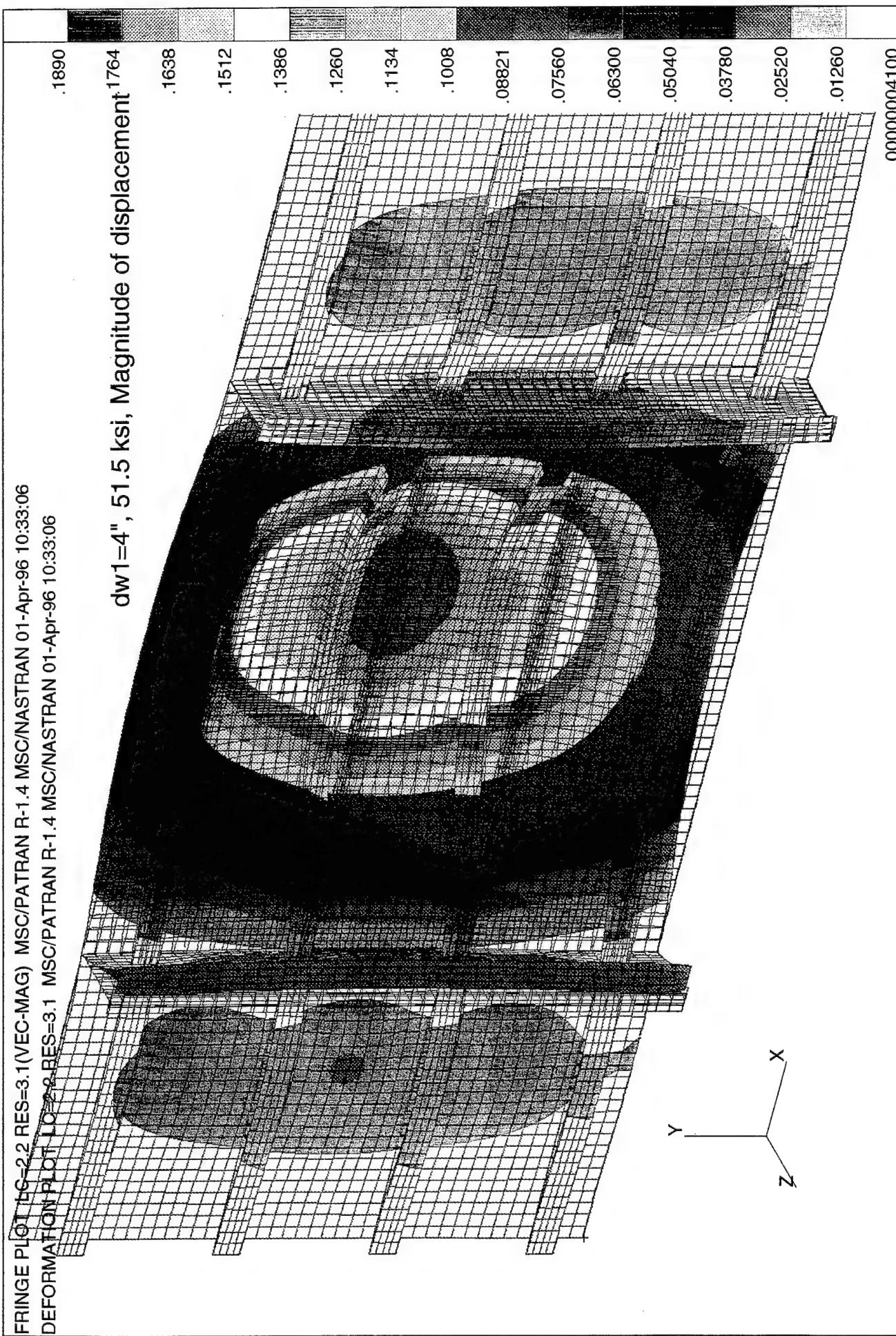


Figure 21

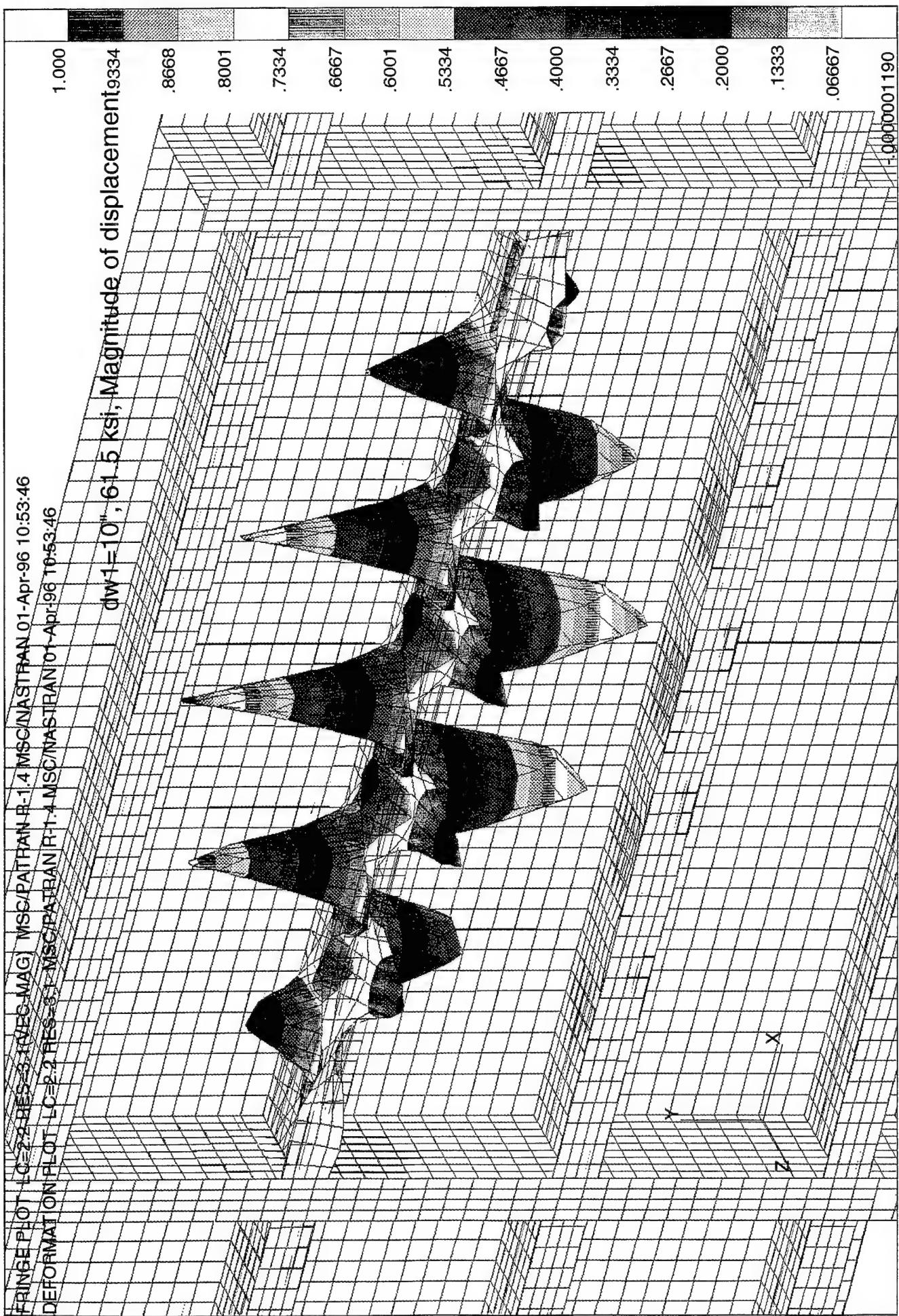


Figure 22

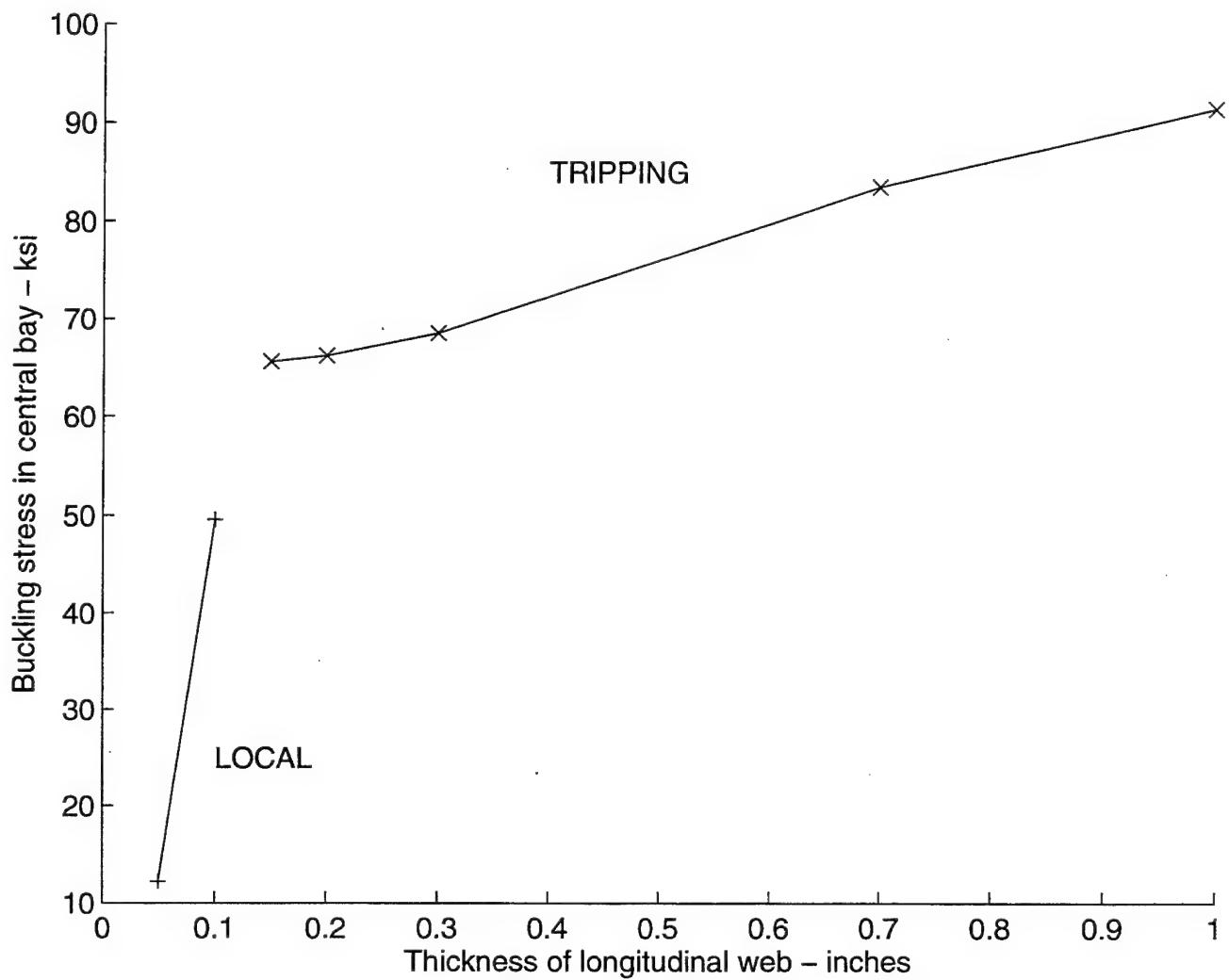


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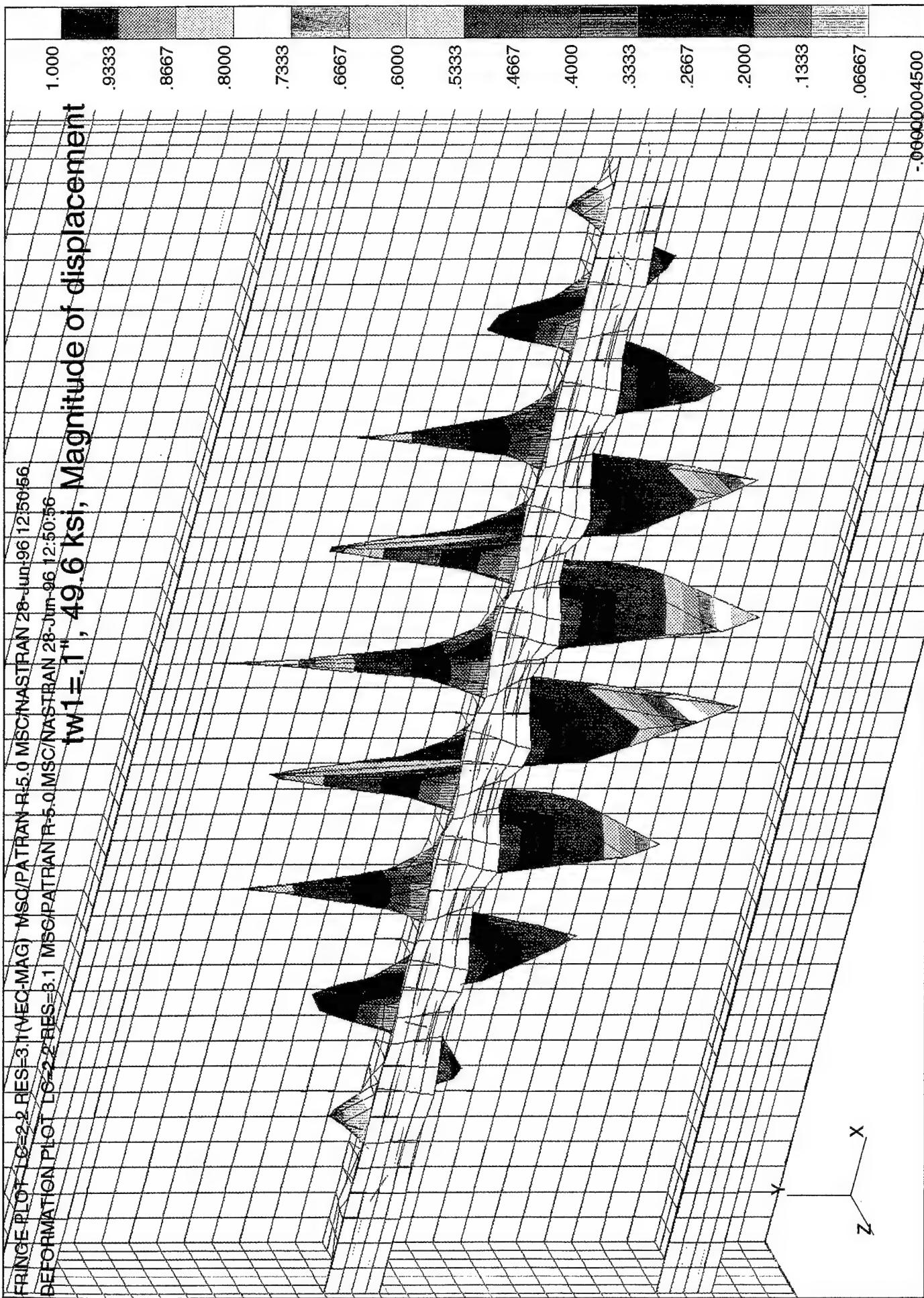


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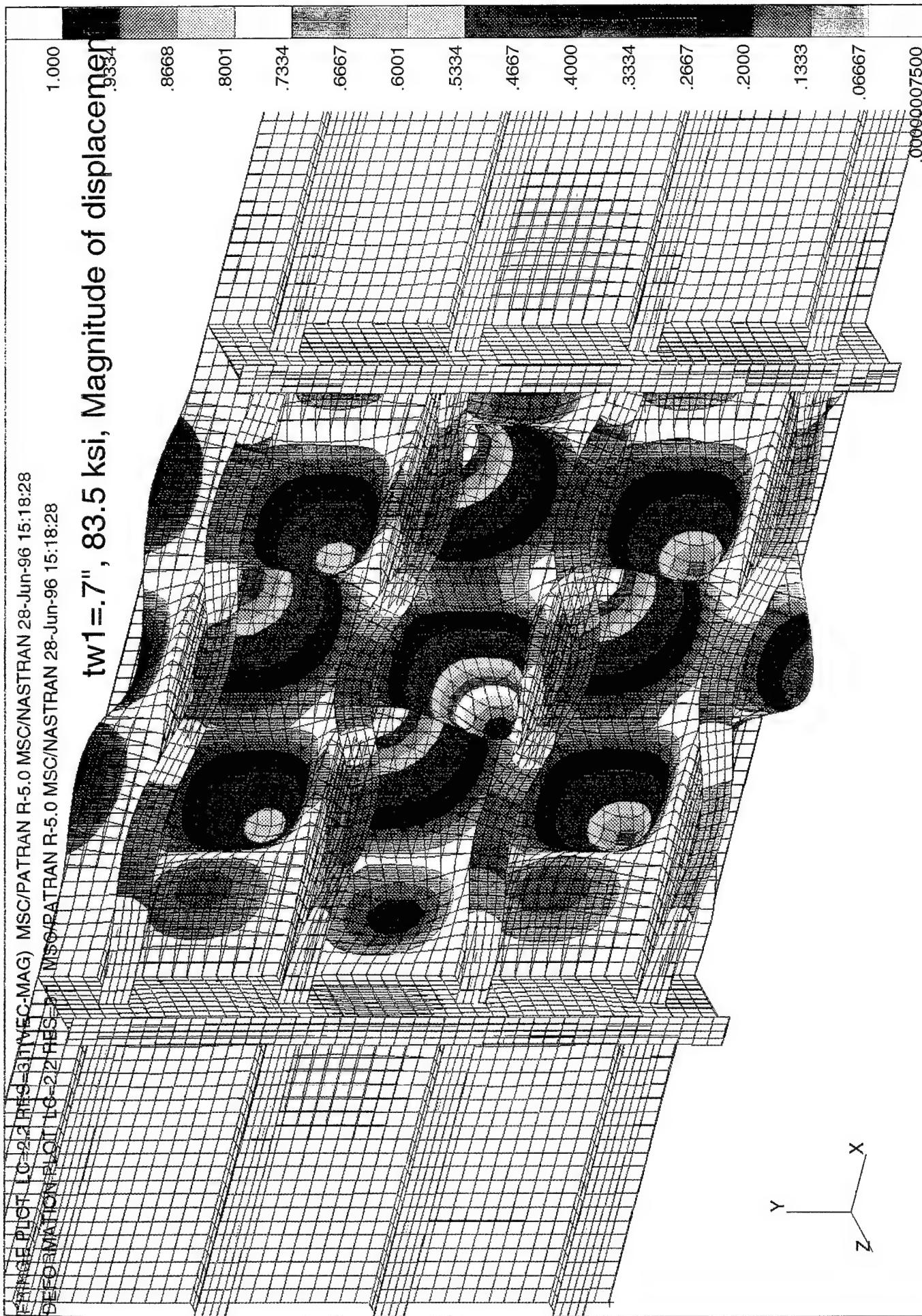


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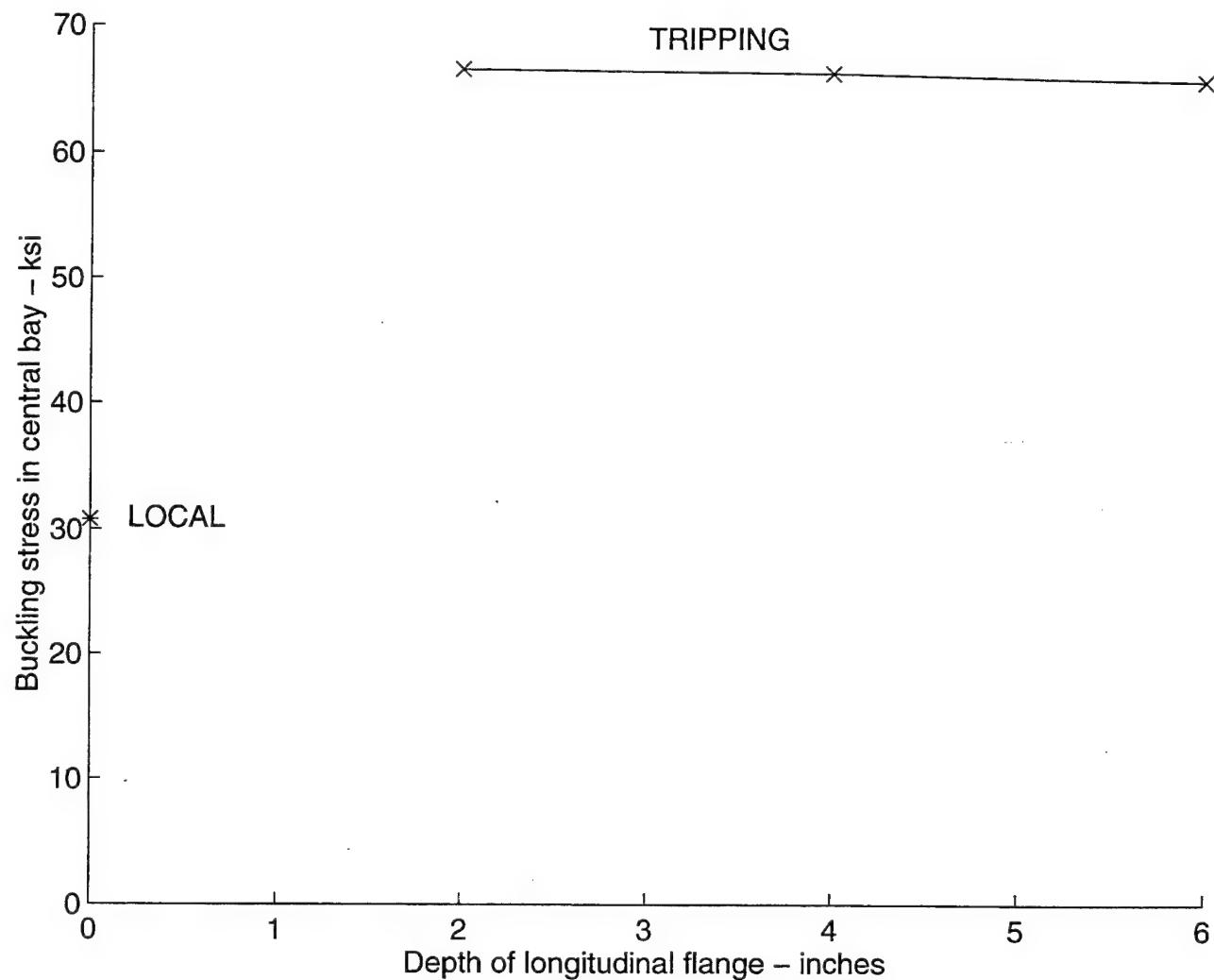


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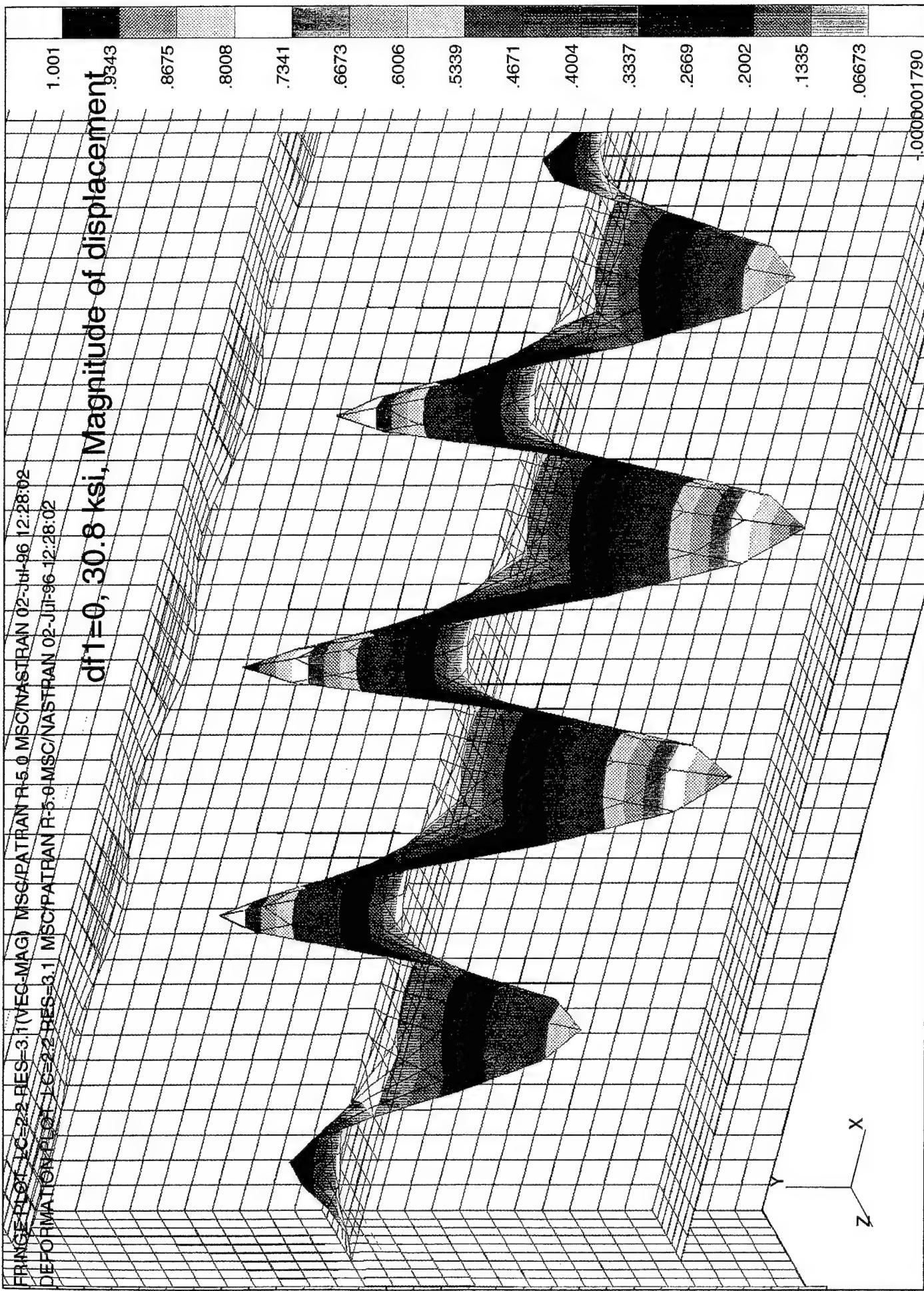


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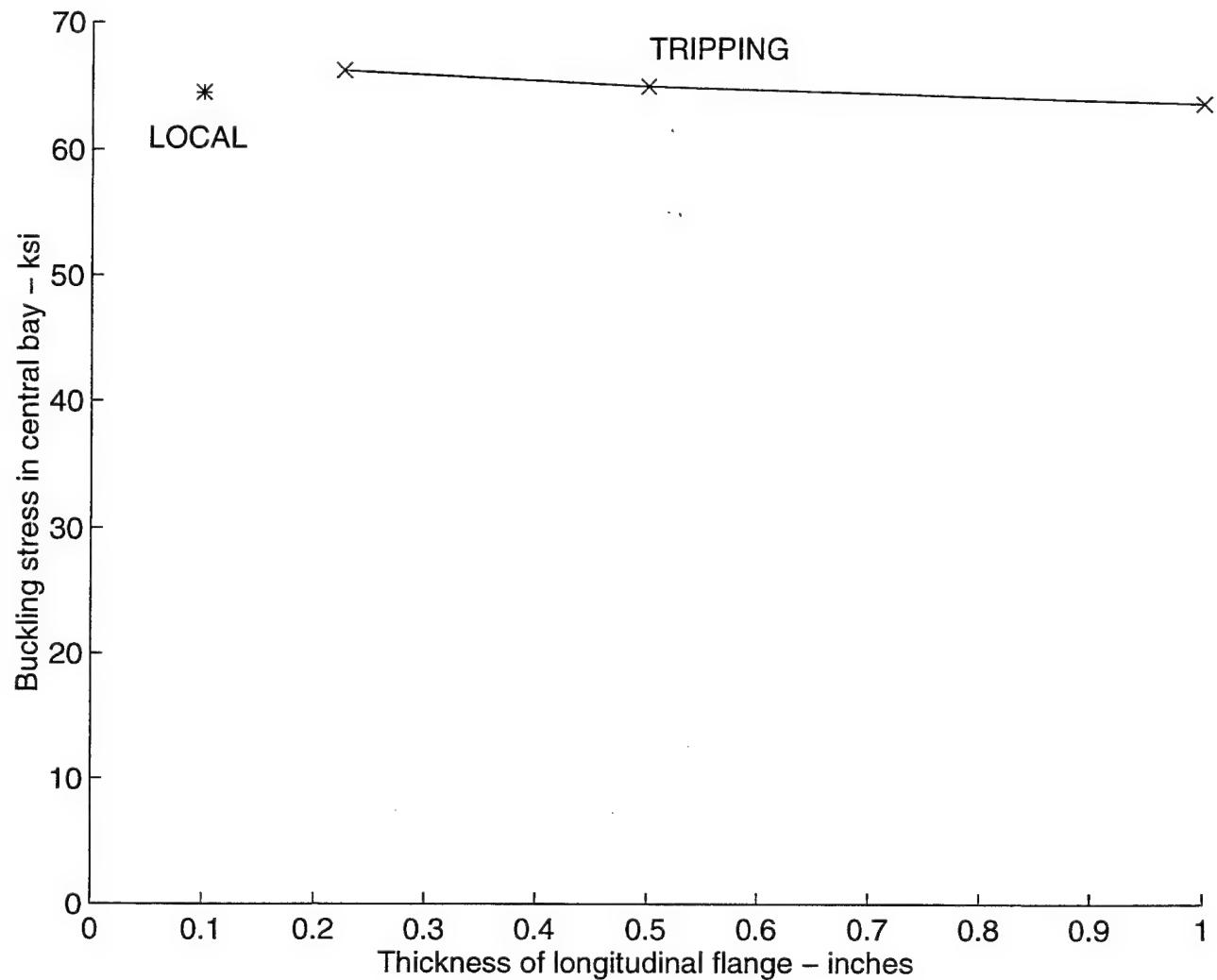


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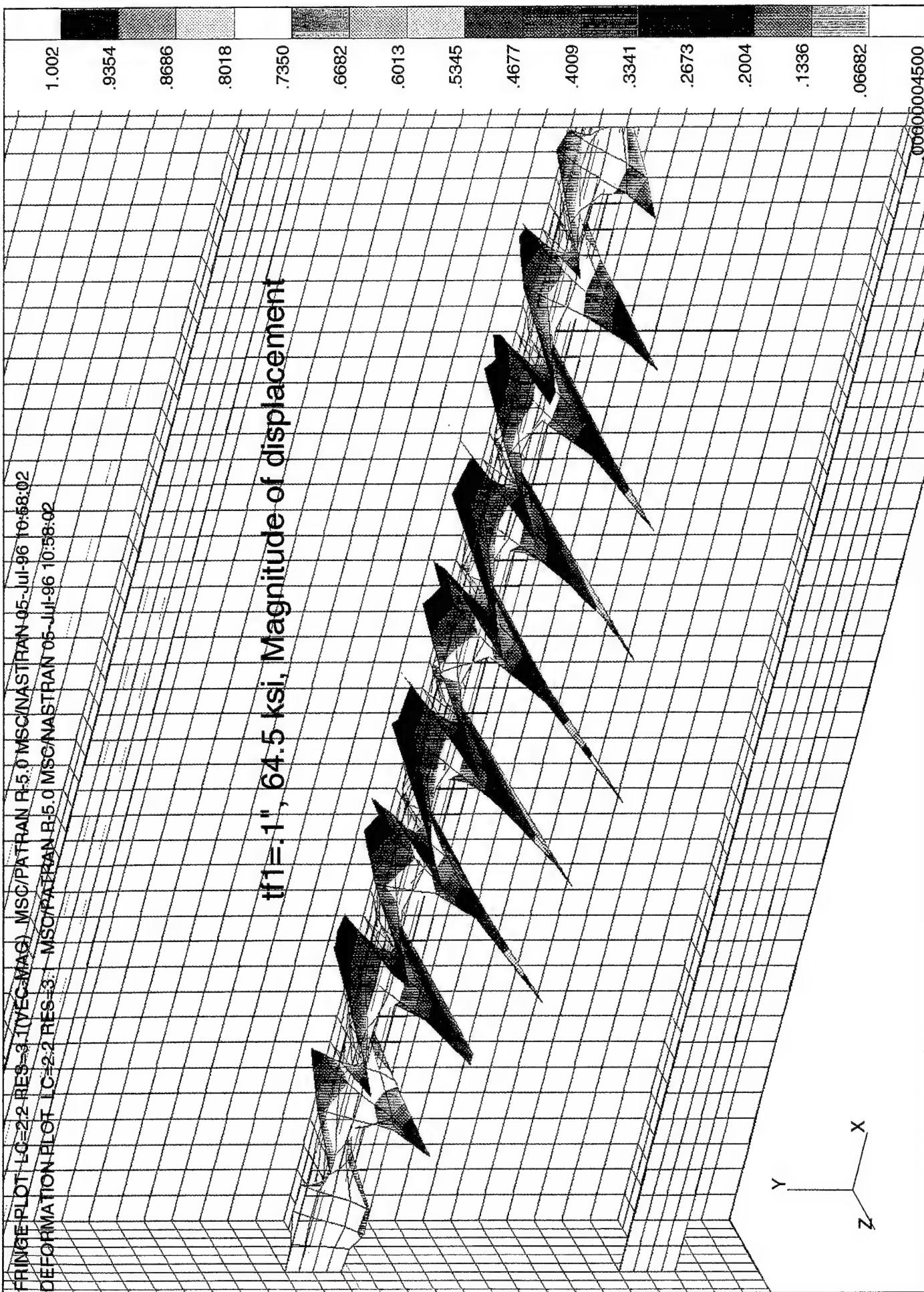


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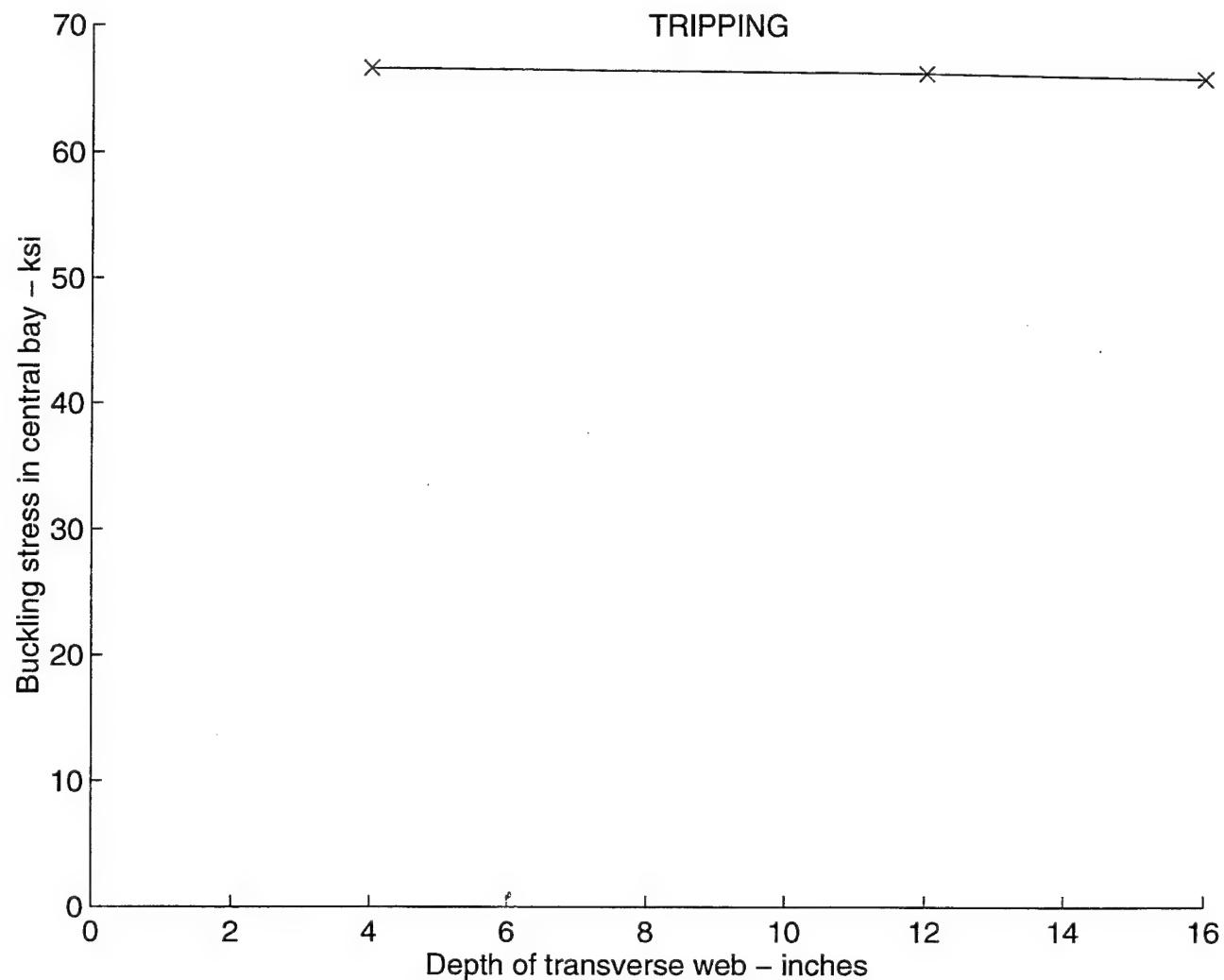


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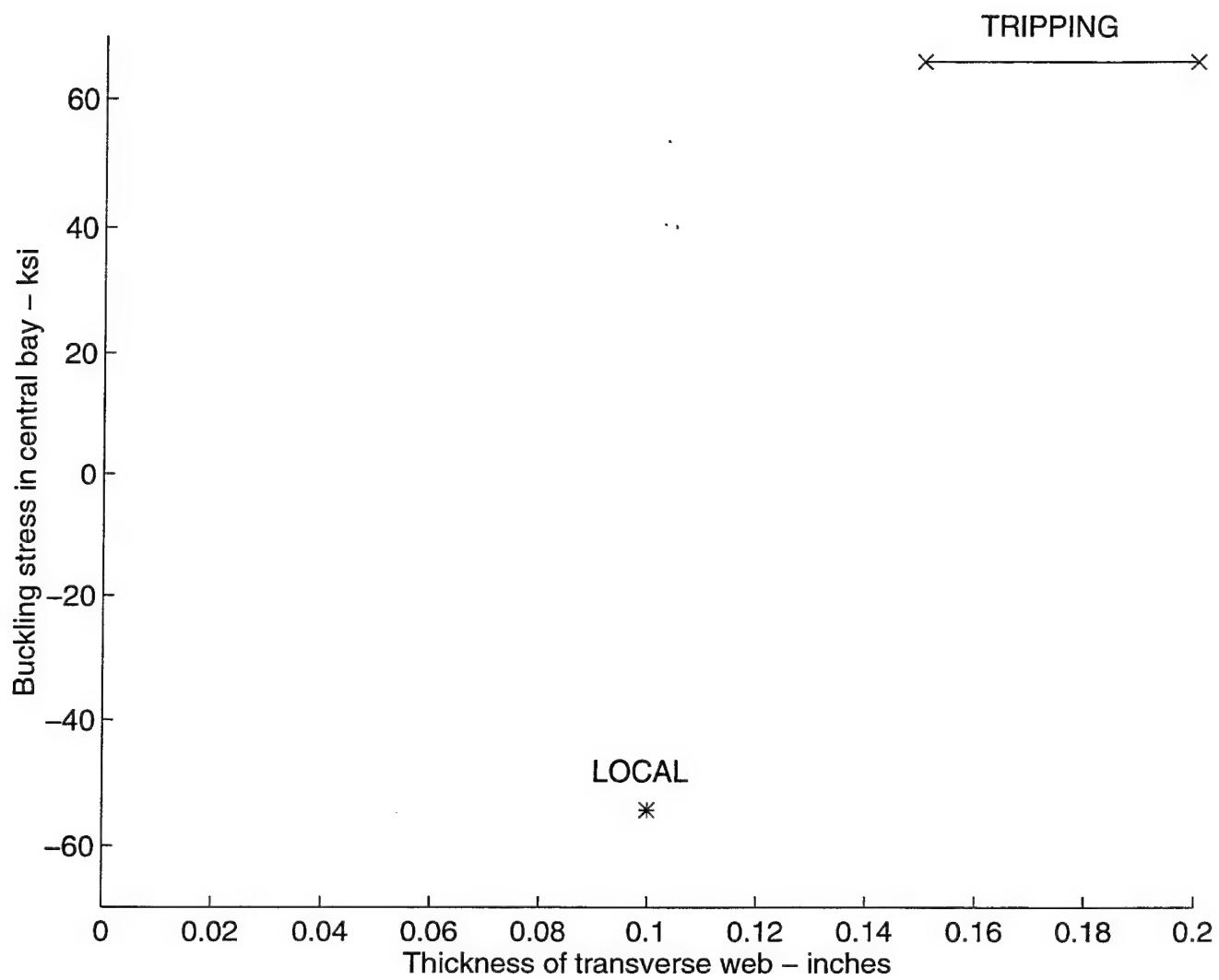


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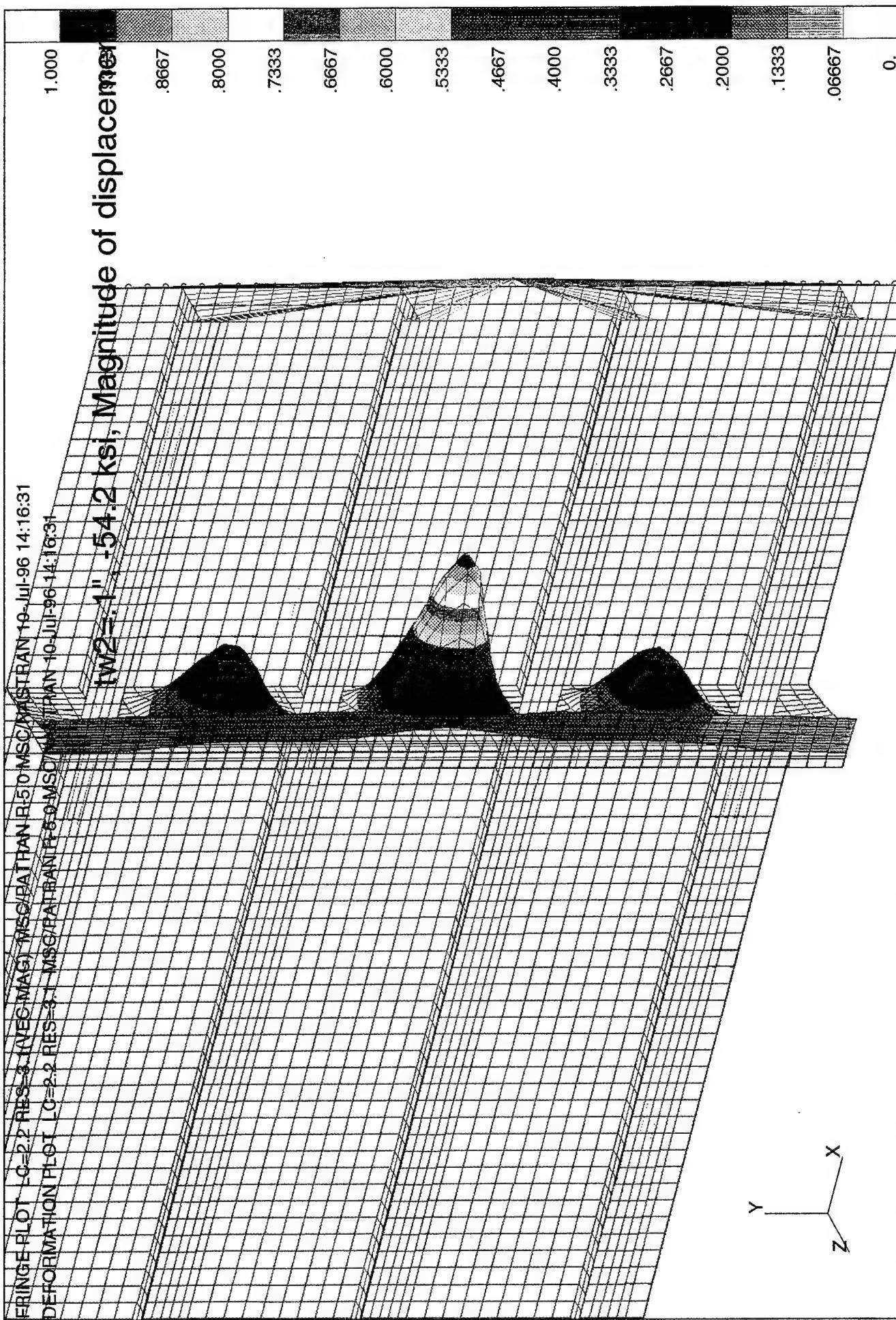


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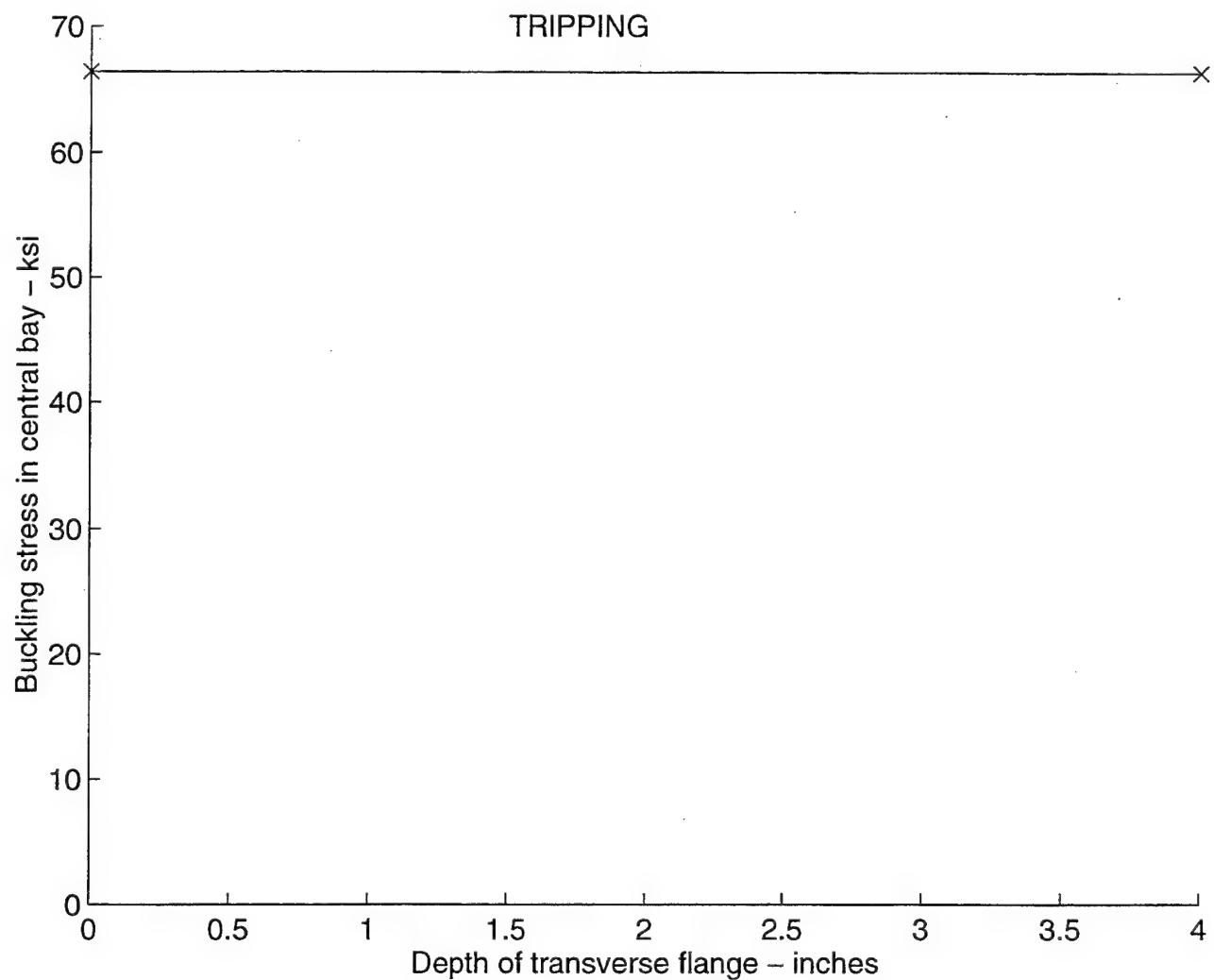


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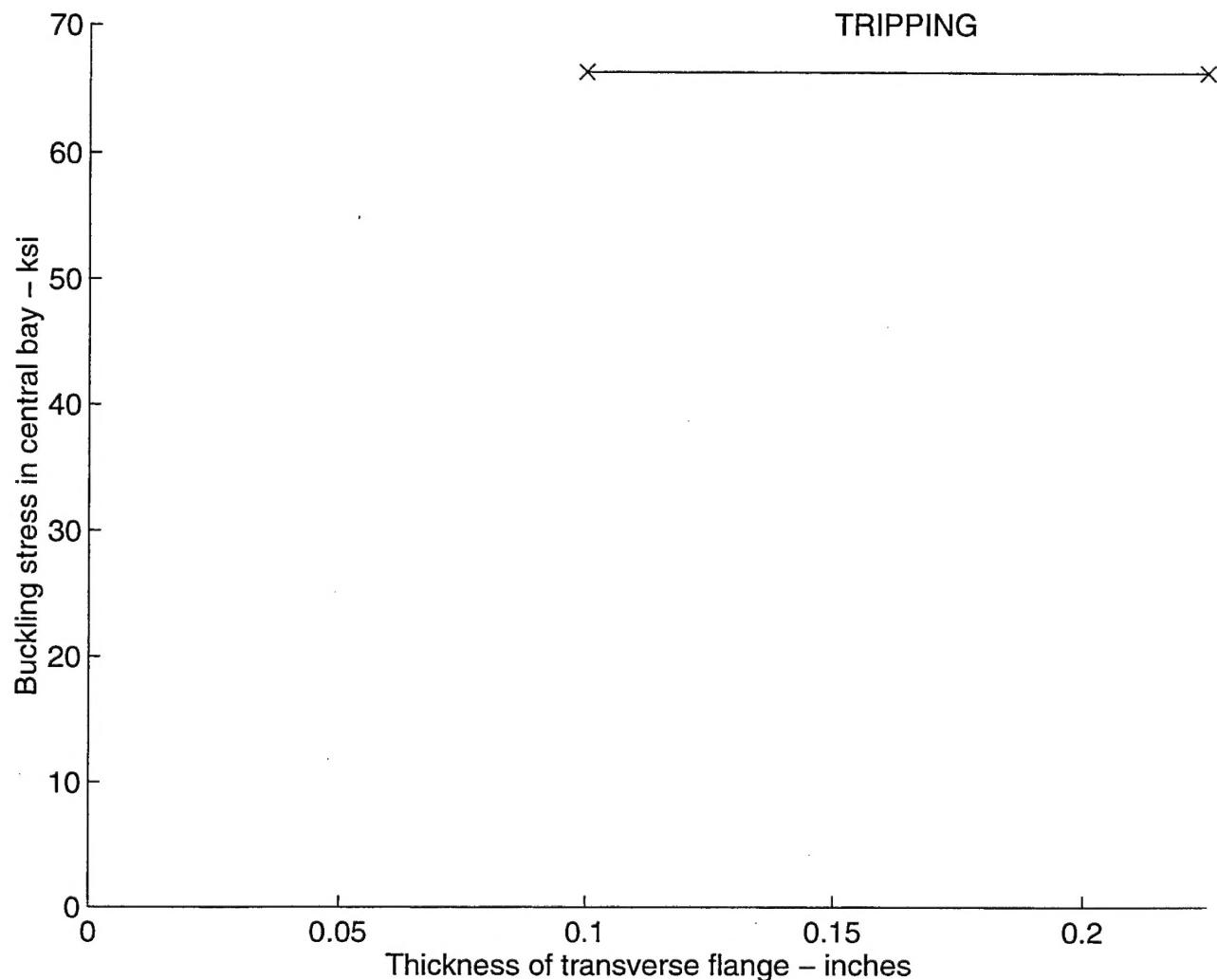


Figure 34

CONCLUSIONS

We can think of the buckling stress as being represented by a surface in n -dimensional parameter space ($n = 11$ in our studies). The surface is continuous but has discontinuous slope at the boundaries between subsurfaces corresponding to TRIPPING, GLOBAL, or LOCAL modes. Our curves of buckling stress versus a parameter are the intersections of the surface with vertical planes passing through a given point on the surface.

To fit the entire surface with an analytical approximation would appear to be a difficult computational problem requiring the calculation of numerous points on the surface. The parameter dependence of the buckling load is coupled, i.e., all the parameters would have to be varied simultaneously.

In future work the finite element eigenvalue code could be used to design the ideal structure. One possible design criterion would be to find the grillage of minimum weight among all possible grillages with buckling stress equal to the yield stress. The candidate points would lie on the intersection of the surface with a horizontal plane.

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